# Domain decomposition preconditioners of Neumann-Neumann type for $h p$-approximations on boundary layer meshes in three dimensions* 

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#### Abstract

We develop and analyze Neumann-Neumann methods for $h p$ finite element approximations of scalar elliptic problems on geometrically refined boundary layer meshes in three dimensions. These are meshes that are highly anisotropic where the aspect ratio typically grows exponentially with the polynomial degree. The condition number of our preconditioners is shown to be independent of the aspect ratio of the mesh and of potentially large jumps of the coefficients. In addition, it only grows polylogarithmically with the polynomial degree, as in the case of $p$ approximations on shape-regular meshes. This work generalizes our previous one on two-dimensional problems in $[40,42]$ and the estimates derived here can be employed to prove condition number bounds for certain types of FETI methods.


Keywords: Domain decomposition, preconditioning, $h p$ finite elements, spectral elements, anisotropic meshes

Subject Classification: 65N22, 65N55, 65N35

[^1]
## 1 Introduction

Solutions of elliptic boundary value problems in polyhedral domains have corner and edge singularities and, in addition, boundary layers may also arise in laminar, viscous, incompressible flows with moderate Reynolds numbers at faces, edges, and corners. Suitably graded meshes, geometrically refined towards corners, edges, and/or faces, can be employed in order to achieve an exponential rate of convergence of $h p$ finite element approximations; see, e.g., $[4,5,26,36,37]$.

Neumann-Neumann (NN) and FETI algorithms are particular iterative substructuring methods and are among the most popular and heavily tested domain decomposition (DD) methods; see, e.g, [19, 13, 24, 7]. Unfortunately, the performance of iterative substructuring methods might be severely compromised if very thin elements and/or subdomains or general non quasiuniform meshes are employed.

Some work has been done on domain decomposition preconditioners for higher order approximations of three-dimensional problems. It is well-known that on shape-regular meshes special care must be taken in the choice of the basis functions in order to produce preconditioners that are robust with respect to the polynomial degree. We mention, e.g., $[21,22,23,29,8,38]$. For $p$ approximations that employ nodal basis functions on Gauss-Lobatto nodes (spectral element approximations), many iterative substructuring methods can be successfully employed and studied; see [33, 34, 30, 32] and the references therein. Some of these ideas can be and have been generalized to $h p$ approximations. We mention, e.g., $[2,1,28,14,20,3,18]$ and the references therein and, in particular, [15] for three-dimensional problems. In all the above-mentioned works, however, the finite element mesh is assumed to be shape-regular and robustness with respect to the aspect ratio is not in general ensured and often unlikely to hold in practice.

In [40, 42], we showed that NN and FETI methods can be successfully devised for the particular geometrically refined boundary layer meshes commonly used for $h p$ finite element approximations of two-dimensional problems. Indeed, these meshes are highly anisotropic, but of a particular type:

1. they are obtained by refining an initial shape-regular mesh (macromesh);
2. refinement is only carried towards the boundary of the computational domain.

These properties, also shared by three-dimensional meshes, allowed us to obtain condition number bounds for the corresponding preconditioned operators that only grow polylogarithmically with the polynomial degree, as is the case of $p$ approximations on shape regular meshes. Our understanding and analysis was confirmed by numerical experiments. In particular, we choose the macromesh
as a decomposition into substructures in such a way that subdomains are shaperegular. Roughly speaking, the reason why such favourable condition numbers are retained lies on the fact that upper bounds come from stable decompositions of finite element functions into components associated to geometrical objects (typically vertices and edges of the subdomains in two dimensions). Because of our particular meshes, only components associated to internal vertices need to be considered, i.e., relative to vertices in a neighborhood of which the mesh is shape-regular.

Three-dimensional boundary layer meshes also share the two characteristics mentioned above. However, stable decompositions now involve face and wirebasket components, where the wirebasket is the union of the subdomain edges and vertices that do not lie on the external boundary of the computational domain. By considering, for instance, an edge $E$ of a macroelement that share a face with $\Omega$ (see the face patch in Figures 1, left, or 2), decoupling of face and wirebasket components is now also performed close to $\partial \Omega$, and thus where the mesh is not shape-regular. In this work, we are however able to provide condition number bounds that only grow polylogarithmically with the polynomial degree, as in the two-dimensional case, and are independent of arbitrarily large aspect ratios of the mesh.

The core of this work lies in the careful modification and derivation of certain Sobolev type inequalities that are independent of the aspect ratio of the mesh for wirebasket and face components of finite element functions; see section 7 . Provided such inequalities are available, the definition of the algorithms and their analysis are fairly standard procedures in DD methods and proceed as in the two-dimensional case in [40]. Here, we will only consider Neumann-Neumann methods, but note that the estimates derived can be employed for the analysis of one-level FETI methods in a straightforward way; see [17, 40].

We limit our analysis to the case of nodal basis functions built on GaussLobatto nodes. In addition, we only consider the model problem (1), which does not have boundary layers but only corner and edge singularities. However, our tensor-product meshes can also be employed when only singularities are present and do not require the use of hanging nodes. We recall that numerical results in [42] for two-dimensional problems showed that better performance is obtained for certain singularly perturbed problems which exhibit boundary layers. In addition, a linear dependence in $k$ for the condition number was observed for problems with geometric refinement towards interfaces that lie in the interior of the computational domain.

The remainder of this paper is organized as follows: in sections 2 and 3, we introduce our continuous and discrete problems, respectively. Geometric boundary layer meshes are introduced in section 4. A particular choice of basis functions is given in section 5 and our Neumann-Neumann preconditioners are defined in section 6 . Section 7 is the core of this work and is devoted to the proof of some discrete Sobolev type inequalities. Comparison results for certain discrete harmonic extensions are given in section 8. Condition number bounds
are then proven in section 9. Section 10 contains some numerical results, while some concluding remarks and perspectives are presented in section 11.

## 2 Problem setting

We consider a linear, elliptic problem on a bounded polyhedral domain $\Omega \subset \mathbb{R}^{3}$ of unit diameter, formulated variationally as:
Find $u \in H_{0}^{1}(\Omega)$, such that

$$
\begin{equation*}
a(u, v)=\int_{\Omega} \rho(\mathbf{x}) \nabla u \cdot \nabla v d \mathbf{x}=f(v), \quad v \in H_{0}^{1}(\Omega) \tag{1}
\end{equation*}
$$

As usual, $H^{1}(\Omega)$ is the space of square summable functions with square summable first derivatives, and $H_{0}^{1}(\Omega)$ its subspace of functions that vanish on $\partial \Omega$. The functional $f(\cdot)$ belongs to the dual space $H^{-1}(\Omega)$. Here $\mathbf{x}=(x, y, z)$ denotes the position vector.

The coefficient $\rho(\mathbf{x})>0$ can be discontinuous, with very different values for different subregions, but we allow it to vary only moderately within each subregion. We will in fact assume that the region is the union of elements (also called subdomains, substructures, or macroelements) $\left\{\Omega_{i}\right\}$. Without decreasing the generality of our results, we will only consider the piecewise constant case:

$$
\rho(\mathbf{x})=\rho_{i}, \quad \mathbf{x} \in \Omega_{i} .
$$

In the case of a region of diameter $H_{i}$, such as the substructure $\Omega_{i}$, we use a norm with different relative weights obtained by a simple dilation argument:

$$
\begin{equation*}
\|u\|_{1, \Omega_{i}}^{2}=|u|_{1, \Omega_{i}}^{2}+\frac{1}{H_{i}^{2}}\|u\|_{0, \Omega_{i}}^{2} . \tag{2}
\end{equation*}
$$

Here, $\|\cdot\|_{0, \Omega_{i}}$ and $|\cdot|_{1, \Omega_{i}}$ denote the norm in $L^{2}\left(\Omega_{i}\right)$ and the seminorm in $H^{1}\left(\Omega_{i}\right)$, respectively. In the following we also employ the space $W^{1, \infty}\left(\Omega_{i}\right)$ of bounded functions with bounded derivatives; see, e.g., [27].

## $3 h p$ finite element approximations

We now specify a particular choice of finite element spaces. Given an affine quadrilateral mesh $\mathcal{T}$ of $\Omega$ and a polynomial degree $k \geq 1$, we consider the following finite element spaces:

$$
\begin{equation*}
X=X^{k}(\Omega ; \mathcal{T}):=\left\{u \in H_{0}^{1}(\Omega) \mid \quad u_{\left.\right|_{K}} \in \mathbb{Q}_{k}(K), K \in \mathcal{T}\right\} \tag{3}
\end{equation*}
$$

Here $\mathbb{Q}_{k}(K)$ is the space of polynomials of maximum degree $k$ in each variable on $K$. In the following, we may drop the reference to $k$, $\Omega$, and/or $\mathcal{T}$ whenever there is no confusion.


Figure 1: Hierarchic structure of a boundary layer mesh, with $\sigma=0.5$ and $n=3$.

In this paper, we always assume that the meshes are regular, i.e, the intersection between neighboring elements is either a vertex, or an edge, or a face that is common to the two elements.

A finite element approximation of (1) consists of finding $u \in X$, such that

$$
\begin{equation*}
a(u, v)=f(v), \quad v \in X \tag{4}
\end{equation*}
$$

## 4 Geometric boundary layer meshes

In order to resolve boundary layers and/or singularities, geometrically graded meshes can be employed. They are determined by a mesh grading factor $\sigma \in$ $(0,1)$ and a refinement level $n \geq 0$. The number of layers is $n+1$ and the thinnest layer has a width proportional to $\sigma^{n}$. Robust exponential convergence of $h p$ finite element approximations is achieved if $n$ is suitably chosen. For singularity resolution, $n$ is required to be proportional to the polynomial degree $k$; see $[4,5]$. For boundary layers, the width of the thinnest layer needs to be comparable to that of the boundary layer; see $[26,36,37]$.

A geometric boundary layer mesh $\mathcal{T}=\mathcal{T}_{b l}^{n, \sigma}$ is, roughly speaking, the tensor product of meshes that are geometrically refined towards the faces. Figure 1 shows the construction of a geometric boundary layer mesh $\mathcal{T}_{b l}^{n, \sigma}$.

The mesh $\mathcal{T}_{b l}^{n, \sigma}$ is built by first considering an initial shape-regular macrotriangulation $\mathcal{T}_{m}$, possibly consisting of just one element, which is successively refined. This process is illustrated in Figure 1. Every macroelement can be refined isotropically (not shown) or anisotropically in order to obtain so-called face, edge, or corner patches (Figure 1, level 2). Here and in the following, we only consider patches obtained by triangulating the reference cube $\hat{Q}:=I^{3}$, with $I:=(-1,1)$. A patch for an element $K_{m} \in \mathcal{T}_{m}$ is obtained by using an
affine mapping $F_{K_{m}}: \hat{Q} \rightarrow K_{m}$. The stability properties proven for patches on the reference cube are equally valid for an arbitrary shape-regular element $K_{m} \in \mathcal{T}_{m}$, with a constant that is independent of the diameter of $K_{m}$.

A face patch is given by an anisotropic triangulation of the form

$$
\begin{equation*}
\mathcal{T}_{f}:=\left\{K_{x} \times I \times I \mid \quad K_{x} \in \mathcal{T}_{x}\right\} \tag{5}
\end{equation*}
$$

where $\mathcal{T}_{x}$ is a mesh of $I$, geometrically refined towards, say $x=1$, with grading factor $\sigma \in(0,1)$ and $n$ levels of refinement; see Figure 1 (level 2, left). We note that the mesh $\mathcal{T}_{x} \times\{I\}$ of $\hat{S}:=I^{2}$ is a two-dimensional edge patch.

An edge patch is given by a triangulation

$$
\begin{equation*}
\mathcal{T}_{e}=\mathcal{T}_{e}^{b l}:=\left\{K_{x} \times K_{y} \times I \mid \quad K_{x} \in \mathcal{T}_{x}, K_{y} \in \mathcal{T}_{y}\right\}=\left\{K_{x y} \times I \mid \quad K_{x y} \in \mathcal{T}_{x y}\right\} \tag{6}
\end{equation*}
$$

where $\mathcal{T}_{x}$ and $\mathcal{T}_{y}$ are meshes of $I$, geometrically refined towards, say $x=1$ and $y=1$, respectively, with grading factor $\sigma \in(0,1)$ and total number of layers $n$; see Figure 1 (level 2, center). The mesh $\mathcal{T}_{x y}$ of $\hat{S}$ is a two-dimensional corner patch.

In a similar way, we can define a corner patch $\mathcal{T}_{c}$ :

$$
\mathcal{T}_{c}=\mathcal{T}_{c}^{b l}:=\left\{K_{x} \times K_{y} \times K_{z} \mid \quad K_{x} \in \mathcal{T}_{x}, K_{y} \in \mathcal{T}_{y}, K_{z} \in \mathcal{T}_{z}\right\}
$$

where $\mathcal{T}_{x}, \mathcal{T}_{y}$, and $\mathcal{T}_{z}$ are meshes of $I$, geometrically refined towards, say $x=1$, $y=1$, and $z=1$, respectively; see Figure 1 (level 2, right).

We note that every element $\hat{K}$ of $\mathcal{T}_{f}, \mathcal{T}_{e}$, and $\mathcal{T}_{c}$ on the reference cube is of the form $\left(0, h_{x}\right) \times\left(0, h_{y}\right) \times\left(0, h_{z}\right)$ (after a possible translation and rotation) and is thus obtained from the reference element by an affine mapping $F_{\hat{K}}: \hat{Q} \rightarrow \hat{K}$ of the form

$$
\left[\begin{array}{lll}
x & y & z
\end{array}\right]^{T}=\left[\begin{array}{lll}
\left(h_{x} / 2\right)(\hat{x}+1) & \left(h_{y} / 2\right)(\hat{y}+1) & \left(h_{z} / 2\right)(\hat{z}+1) \tag{7}
\end{array}\right]^{T} .
$$

The aspect ratio of $\hat{K}$ is the maximum of all possible ratios of $h_{x}, h_{y}$, and $h_{z}$. Since the macromesh consists of affinely mapped elements $K_{m}$, every element $K$ of the global mesh $\mathcal{T}=\mathcal{T}_{b l}^{n, \sigma}$ is obtained from the reference element by combining two affine mappings

$$
\begin{equation*}
K=F_{K}(\hat{Q})=F_{K_{m}}\left(F_{\hat{K}}(\hat{Q})\right), \quad K \subset K_{m} \in \mathcal{T}_{m} \tag{8}
\end{equation*}
$$

Since $\mathcal{T}_{m}$ is shape-regular, the aspect ratio is determined only by $F_{\hat{K}}$; cf. (7). Finally we note that the aspect ratio of the mesh is determined by $\sigma$ and $n$, and is proportional to $\sigma^{-n}$.

As in [40], our analysis will be made for a prototype mesh, obtained from a shape-regular (not necessarily quasi-uniform) macromesh, by refining elements that only touch $\partial \Omega$, either as corner, edge, or face patches. Such meshes only consist of four types of patches: unrefined, face, edge, and corner patches. We also recall that in practical applications $\sigma$ is bounded away from one and zero.

## 5 Basis functions on Gauss-Lobatto nodes

For the space $X^{k}(\Omega ; \mathcal{T})$, we choose nodal basis functions on the Gauss-Lobatto nodes. We denote by $G L L(k)$ the set of Gauss-Lobatto points $\left\{\xi_{i} ; 0 \leq i \leq k\right\}$ on $I=(-1,1)$ in increasing order and by $\left\{w_{i}>0\right\}$ the corresponding weights; see [6, Sect. 4]. We recall that the quadrature formula based on $G L L(k)$ has order $2 k-1$ and, in addition,

$$
\begin{equation*}
\|u\|_{0, I}^{2} \leq \sum_{i=0}^{k} u\left(\xi_{i}\right)^{2} w_{i} \leq 3\|u\|_{0, I}^{2}, \quad u \in \mathbb{Q}_{k}(I) \tag{9}
\end{equation*}
$$

see [6, Rem. 13.3].
For the reference cube $\hat{Q}=(-1,1)^{3}$ we set $G L L(k)^{3}=\left\{\xi_{i j l}=\left(\xi_{i}, \xi_{j}, \xi_{l}\right) ; 0 \leq\right.$ $i, j, l \leq k\}$. In the following, we use the same notation for the mapped GaussLobatto nodes and corresponding weights for an affinely mapped element $K \in$ $\mathcal{T}$.

Given the nodes $G L L(k)^{3}$, our basis functions on $\mathbb{Q}_{k}(\hat{Q})$ are the tensor product of $k$-th order Lagrange interpolating polynomials on $G L L(k)$, defined by

$$
\begin{equation*}
\hat{l}_{i}\left(\xi_{j}\right)=\delta_{i j} \tag{10}
\end{equation*}
$$

On the reference element we can write

$$
\begin{equation*}
u(x, y, z)=\sum_{i=0}^{k} \sum_{j=0}^{k} \sum_{l=0}^{k} u\left(\xi_{i}, \xi_{j}, \xi_{l}\right) \hat{l}_{i}(x) \hat{l}_{j}(y) \hat{l}_{l}(z), \quad u \in \mathbb{Q}_{k}(\hat{Q}) \tag{11}
\end{equation*}
$$

For a general element in $\mathcal{T}$, basis functions are obtained by mapping those on the reference element. Interior local basis functions correspond to GLL nodes inside $\hat{Q}$ (all local indices differ from 0 and $k$ ).

Equation (11) defines an interpolation operator $I^{k}$ on the reference element

$$
I^{k} u(x, y, z):=\sum_{i=0}^{k} \sum_{j=0}^{k} \sum_{l=0}^{k} u\left(\xi_{i}, \xi_{j}, \xi_{l}\right) \hat{l}_{i}(x) \hat{l}_{j}(y) \hat{l}_{l}(z)
$$

The points $G L L(k)^{3}$ define a triangulation $\mathcal{T}_{k}=\mathcal{T}_{k}(\hat{Q})$ of $\hat{Q}$ in a natural way, consisting of $k^{3}$ parallelepipeds. Let $Y^{h}=Y^{h}(\hat{Q})=X^{1}\left(\hat{Q} ; \mathcal{T}_{k}\right)$ be the space of piecewise trilinear functions on this mesh. We also denote $Y^{k}=Y^{k}(\hat{Q})=$ $\mathbb{Q}_{k}(\hat{Q})$. The aspect ratio of $\mathcal{T}_{k}$ is of the order of $k$; see [10, Pg. 27] for details. In a similar way we can consider a Gauss-Lobatto mesh on an affinely mapped element $K$ by simply mapping the GLL mesh on $\hat{Q}$. In the following, we will use the notations $\mathcal{T}_{k}=\mathcal{T}_{k}(K), Y^{h}=Y^{h}(K)$, and $Y^{k}=Y^{k}(K)$, to denote the GLL mesh, the piecewise trilinear finite element space, and $\mathbb{Q}_{k}$, respectively, for a mapped element. If the aspect ratio of $K$ is, e.g., $h_{x} / h_{y}$ (cf. (7) and (8)), then that of the corresponding $\mathcal{T}_{k}$ is $\left(h_{x} / h_{y}\right) k$.

There is a one-to-one correspondence between $Y^{h}$ and $Y^{k}$ given by

$$
I^{k}: Y^{h} \rightarrow Y^{k}, \quad I^{h}: Y^{k} \rightarrow Y^{h}
$$

where $I^{h}$ is the nodal interpolation operator on $Y^{h}$. We use the notation $u_{h} \in$ $Y^{h}$ and $u_{k} \in Y^{k}$ in order to denote two corresponding functions.
Lemma 5.1 Let $\hat{K}=\left(0, h_{x}\right) \times\left(0, h_{y}\right) \times\left(0, h_{z}\right)$. Then there exist positive constants $c$ and $C$, such that, for $u_{h} \in Y^{h}(\hat{K})$,

$$
\begin{aligned}
& c\left\|u_{h}\right\|_{0, \hat{K}} \leq\left\|u_{k}\right\|_{0, \hat{K}} \leq C\left\|u_{h}\right\|_{0, \hat{K}} \\
& c\left\|\partial_{x}\left(u_{h}\right)\right\|_{0, \hat{K}} \leq\left\|\partial_{x}\left(u_{k}\right)\right\|_{0, \hat{K}} \leq C\left\|\partial_{x}\left(u_{h}\right)\right\|_{0, \hat{K}}
\end{aligned}
$$

with, in particular, $c$ and $C$ independent of $h_{x}, h_{y}, h_{z}$, and $k$. Similar bounds hold for the $y$ and $z$ derivatives. If $K \in \mathcal{T}$ is given by (8), then, for $u_{h} \in Y^{h}(K)$,

$$
\begin{aligned}
& c\left\|u_{h}\right\|_{0, K} \leq\left\|u_{k}\right\|_{0, K} \leq C\left\|u_{h}\right\|_{0, K} \\
& c\left|u_{h}\right|_{1, K} \leq\left|u_{k}\right|_{1, K} \leq C\left|u_{h}\right|_{1, K}
\end{aligned}
$$

where the constants are independent of the diameter and the aspect ratio of $K$, and $k$.

The proof of the above result can be found in [9, Sect. 2] for $K=\hat{Q}$. For an affinely mapped element a scaling argument can be used. We note that thanks to Lemma 5.1 we can equivalently work with functions in $Y^{k}$ or $Y^{h}$.

The following result can be found in [10, Lem. 3.3.3].
Lemma 5.2 Let $\hat{K}=\left(0, h_{x}\right) \times\left(0, h_{y}\right) \times\left(0, h_{z}\right)$ and $u_{h} \in Y^{h}(\hat{K})$. Given $\theta \in$ $W^{1, \infty}(\hat{K})$, with

$$
\|\theta\|_{\infty, \hat{K}} \leq C, \quad\|\nabla \theta\|_{\infty, \hat{K}} \leq C / r
$$

then

$$
\begin{aligned}
& \left\|I^{h}\left(\theta u_{h}\right)\right\|_{0, \hat{K}}^{2} \leq C\left\|u_{h}\right\|_{0, \hat{K}}^{2} \\
& \left\|\partial_{x} I^{h}\left(\theta u_{h}\right)\right\|_{0, \hat{K}}^{2} \leq C\left(\left|u_{h}\right|_{1, \hat{K}}^{2}+r^{-2}\left\|u_{h}\right\|_{0, \hat{K}}^{2}\right),
\end{aligned}
$$

where $C$ is independent of $h_{x}, h_{y}, h_{z}$, and $k$. Similar bounds hold for the $y$ and $z$ derivatives. If $K \in \mathcal{T}$ is given by (8), then, for $u_{h} \in Y^{h}(K)$,

$$
\begin{aligned}
& \left\|I^{h}\left(\theta u_{h}\right)\right\|_{0, K}^{2} \leq C\left\|u_{h}\right\|_{0, K}^{2} \\
& \left|I^{h}\left(\theta u_{h}\right)\right|_{1, K}^{2} \leq C\left(\left|u_{h}\right|_{1, K}^{2}+r^{-2}\left\|u_{h}\right\|_{0, K}^{2}\right)
\end{aligned}
$$

where $C$ is independent of the diameter and the aspect ratio of $K$, and $k$.
Given an element $\hat{K}=\left(0, h_{x}\right) \times\left(0, h_{y}\right) \times\left(0, h_{z}\right)$ and a coordinate direction, say $x$, , let $a, b, c$, and $d$ be the vertices of a face of $\hat{K}$ perpendicular to this direction, and let $a^{\prime}, b^{\prime}, c^{\prime}$, and $d^{\prime}$ be the corresponding points on the parallel face. The following lemma relies on trivial properties of trilinear functions; cf. [10, Lem. 3.3.1].

Lemma 5.3 Let $\hat{K}=\left(0, h_{x}\right) \times\left(0, h_{y}\right) \times\left(0, h_{z}\right)$ and $a, b$, $c$, and $d$ be the vertices of a face of $\hat{K}$ perpendicular to the $x$ direction. Then there are constants independent of $h_{x}, h_{y}$, and $h_{z}$, such that, if $u$ is trilinear on $\hat{K}$,

$$
\begin{aligned}
& c\|u\|_{0, \hat{K}}^{2} \leq h_{x} h_{y} h_{z} \sum_{\mathbf{x}=a, b, c, d}\left(u(\mathbf{x})^{2}+u\left(\mathbf{x}^{\prime}\right)^{2}\right) \leq C\|u\|_{0, \hat{K}}^{2} \\
& c\left\|\partial_{x} u\right\|_{0, \hat{K}}^{2} \leq\left(h_{x} h_{y} h_{z} / h_{x}^{2}\right) \sum_{\mathbf{x}=a, b, c, d}\left(u(\mathbf{x})-u\left(\mathbf{x}^{\prime}\right)\right)^{2} \leq C\left\|\partial_{x} u\right\|_{0, \hat{K}}^{2}, \\
& c\left\|\partial_{x} u\right\|_{\infty, \hat{K}}^{2} \leq h_{x}^{-2} \sum_{\mathbf{x}=a, b, c, d}\left(u(\mathbf{x})-u\left(\mathbf{x}^{\prime}\right)\right)^{2} \leq C\left\|\partial_{x} u\right\|_{\infty, \hat{K}}^{2} .
\end{aligned}
$$

Similar bounds hold for the $y$ and $z$ derivatives.

## 6 Neumann-Neumann methods

Iterative substructuring methods rely on a non-overlapping partition into substructures. We mention [39, Ch. 4] as a general reference to this section. In our algorithms the substructures are chosen as the macroelements in $\mathcal{T}_{m}=\left\{\Omega_{i} \mid 1 \leq\right.$ $i \leq N\}$. We recall that the macroelements are shape-regular. This appears to be essential for the analysis and good performance.

We define the boundaries $\Gamma_{i}=\partial \Omega_{i} \backslash \partial \Omega$ and the interface $\Gamma$ as their union. We remark that $\Gamma$ is the union of the interior subdomain faces, regarded as open sets, which are shared by two subregions, and subdomain edges and vertices, which are shared by more than two subregions. Vertices can only be endpoints of edges. In the following, we tacitly assume that points on $\partial \Omega$ are excluded from the geometrical objects that we consider, or, in other words, we will only deal with geometrical objects (faces, edges, vertices, ...) that belong to $\Gamma$. We denote the faces of $\Omega_{i}$ by $F^{i j}$, its edges by $E^{i j}$, its vertices by $V^{i j}$, and its wirebasket, defined as the union of its edges and vertices, by $W^{i}$. Occasionally, we will also use faces, edges, and vertices with one or no superscript. If a vertex (edge) lies on $\partial \Omega$ we will regard it as part of the internal edge (resp., face) that shares it with $\partial \Omega$.

When restricted to the subdomain $\Omega_{i}$, the global triangulation $\mathcal{T}$ determines a local mesh $\mathcal{T}_{i}$. This mesh can be of four types: face, edge, corner, or consisting of just one element. We define the local spaces $X_{i}=X^{k}\left(\Omega_{i} ; \mathcal{T}_{i}\right)$, of local finite element functions that vanish on $\partial \Omega \cap \partial \Omega_{i}$

In our analysis, we will also employ the GLL mesh $\mathcal{T}_{k}\left(\Omega_{i}\right)$ on $\Omega_{i}$, generated by the local GLL meshes $\mathcal{T}_{k}(K)$ for $K \in \mathcal{T}_{i}$. The corresponding space of piecewise trilinear functions on $\mathcal{T}_{k}\left(\Omega_{i}\right)$ that vanish on $\partial \Omega \cap \partial \Omega_{i}$ is denoted by $Y^{h}\left(\Omega_{i}\right)$. We set $Y^{k}\left(\Omega_{i}\right)=X^{k}\left(\Omega_{i} ; \mathcal{T}_{i}\right)$.

We next define the local bilinear forms

$$
a_{i}(u, v)=\int_{\Omega_{i}} \rho_{i} \nabla u \cdot \nabla v d \mathbf{x}, \quad u, v \in X_{i} .
$$

We note that if $\Omega_{i}$ is a floating subdomain (i.e., its boundary does not touch $\partial \Omega), a_{i}(\cdot, \cdot)$ is only positive semi-definite and for $u \in X_{i}$ we have

$$
a_{i}(u, u)=0 \quad \text { iff } \quad u \text { constant in } \Omega_{i} .
$$

The sets of nodal points on $\Gamma_{i}, \Gamma, F^{i j}, E^{i j}$, and $W^{i}$ are denoted by $\Gamma_{i, h}, \Gamma_{h}$, $F_{h}^{i j}, E_{h}^{i j}$, and $W_{h}^{i}$, respectively. We will identify these sets with the corresponding sets of degrees of freedom. As for the corresponding regions, we will also use notations with one or no superscript.

We introduce some spaces defined on the interfaces: $U_{i}$ is the space of restrictions to $\Gamma_{i}$ of functions in $X^{k}\left(\Omega_{i} ; \mathcal{T}_{i}\right)$ and $U$ of restrictions to $\Gamma$ of functions in $X^{k}(\Omega ; \mathcal{T})$. We note that functions in $U_{i}$ and $U$ are uniquely determined by the nodal values in $\Gamma_{i, h}$ and $\Gamma_{h}$, respectively. For every substructure $\Omega_{i}$, there is a natural interpolation operator

$$
R_{i}^{T}: U_{i} \longrightarrow U
$$

that extends a function on $\Gamma_{i}$ to a global function on $\Gamma$ with vanishing degrees of freedom in $\Gamma_{h} \backslash \Gamma_{i, h}$. Its transpose with respect to the Euclidean scalar product $R_{i}: U \rightarrow U_{i}$ extracts the degrees of freedom in $\Gamma_{i, h}$.

Once a vector $u \in X^{k}(\Omega ; \mathcal{T})$ is expanded using the basis functions introduced in section 5, Problem (4) can be written as a linear system

$$
A u=f
$$

We recall that the condition number of $A$ is expected to grow at least as $k^{3} /\left(h_{\text {min }}\right)^{2} \sim k^{3} \sigma^{-2 n} \sim k^{3} \sigma^{-2 k}$ (see [25] for a result in two dimensions) and may thus be extremely large for large values of $k$.

The contributions to the stiffness matrix and the right hand side can be formed one subdomain at a time. The stiffness matrix is then obtained by subassembly of these parts. We will order the nodal points interior to the subdomains first, followed by those on the interface $\Gamma$. Similarly, for the stiffness matrix relative to a substructure $\Omega_{i}$, we have

$$
A^{(i)}=\left(\begin{array}{ll}
A_{I I}^{(i)} & A_{I \Gamma}^{(i)} \\
A_{\Gamma I}^{(i)} & A_{\Gamma \Gamma}^{(i)}
\end{array}\right) .
$$

In a first step of many iterative substructuring algorithms, the unknowns in the interior of the subdomains are eliminated by block gaussian elimination. In this step, the Schur complements, with respect to the variables associated with the boundaries of the individual substructures, are calculated. The resulting linear system can be written as

$$
\begin{equation*}
S u_{\Gamma}=g_{\Gamma} . \tag{12}
\end{equation*}
$$

Given the local Schur complements

$$
S_{i}=A_{\Gamma \Gamma}^{(i)}-A_{I \Gamma}^{(i)^{T}} A_{I I}^{(i)^{-1}} A_{I \Gamma}^{(i)}: U_{i} \longrightarrow U_{i},
$$

we have

$$
S=\sum_{i=1}^{N} R_{i}^{T} S_{i} R_{i}: U \longrightarrow U
$$

and an analogous formula can be found for $g_{\Gamma}$; see [39, Ch. 4].
A function $u^{(i)}$ defined on $\Omega_{i}$ is said to be discrete harmonic on $\Omega_{i}$ if

$$
A_{I I}^{(i)} u_{I}^{(i)}+A_{I \Gamma}^{(i)} u_{\Gamma}^{(i)}=0 .
$$

In this case, it is easy to see that $\mathcal{H}_{i}\left(u_{\Gamma}^{(i)}\right):=u^{(i)}$ is completely defined by its value on $\Gamma_{i}$. The space of piecewise discrete harmonic functions $u$ consists of functions in $X$ that are discrete harmonic on each substructure. In this case, $u=: \mathcal{H}\left(u_{\Gamma}\right)$ is completely defined by its value on $\Gamma$.

Our preconditioners will be defined with respect to the inner product

$$
s(u, v)=u^{T} S v, \quad u, v \in U
$$

It follows immediately from the definition of $S$ that $s(\cdot, \cdot)$ is symmetric and coercive.

The following lemma results from elementary variational arguments.
Lemma 6.1 Let $u_{\Gamma}^{(i)}$ be the restriction of a finite element function to $\Gamma_{i}$. Then the discrete harmonic extension $u^{(i)}=\mathcal{H}_{i}\left(u_{\Gamma}^{(i)}\right)$ of $u_{\Gamma}^{(i)}$ into $\Omega_{i}$ satisfies

$$
a_{i}\left(u^{(i)}, u^{(i)}\right)=\min _{\left.v^{(i)}\right|_{\partial \Omega_{i}}=u_{\Gamma}^{(i)}} a_{i}\left(v^{(i)}, v^{(i)}\right)=u_{\Gamma}^{(i)^{T}} S^{(i)} u_{\Gamma}^{(i)} .
$$

Analogously, if $u_{\Gamma}$ is the restriction of a finite element function to $\Gamma$, the piecewise discrete harmonic extension $u=\mathcal{H}\left(u_{\Gamma}\right)$ of $u_{\Gamma}$ into the interior of the subdomains satisfies

$$
a(u, u)=\min _{\left.v\right|_{\Gamma}=u_{\Gamma}} a(v, v)=s(u, u)=u_{\Gamma}^{T} S u_{\Gamma}
$$

This lemma ensures that instead of working with functions defined on the interface $\Gamma$, we can equivalently work with the corresponding discrete harmonic extensions. For this reason, in the following we will identify spaces of traces on the interfaces, $U_{i}$ and $U$, with spaces of discrete harmonic extensions. We point out however that due to the particular meshes considered, we cannot equivalently work with norms of local discrete harmonic extensions and traces on the subdomain boundaries since our local meshes are not in general quasi-uniform or shape-regular, and stable discrete harmonic extensions cannot be found in general.

Neumann-Neumann methods provide preconditioners for the Schur complement system: instead of solving (12) using, e.g., the conjugate gradient method, they employ an equivalent system involving a preconditioned operator of the form

$$
\hat{S}^{-1} S=P_{N N}=P_{0}+\left(I-P_{0}\right)\left(\sum_{i=1}^{N} P_{i}\right)\left(I-P_{0}\right)
$$

We refer to $[12,24,30,17]$ for some NN methods for the $h$ and $p$ finite element approximations. We are unaware on any such method for $h p$ approximations.

The operators $P_{i}$ are projection-like operators associated to a family of subspaces $U_{i}$ and determined by a set of local bilinear forms defined on them

$$
\tilde{s}_{i}(u, v), \quad u, v \in U_{i}
$$

Given the interpolation operators $R_{i}^{T}: U_{i} \rightarrow U$, we have

$$
\begin{equation*}
P_{i}=R_{i}^{T} \tilde{P}_{i}, \quad \tilde{P}_{i}: U \longrightarrow U_{i} \tag{13}
\end{equation*}
$$

with

$$
\begin{equation*}
\tilde{s}_{i}\left(\tilde{P}_{i} u, v_{i}\right)=s\left(u, R_{i}^{T} v_{i}\right), \quad v_{i} \in U_{i} \tag{14}
\end{equation*}
$$

While $P_{0}$ is associated to a low-dimensional global problem, the others are associated to the single substructures. The remainder of this section is devoted to the definition of the various components of $P_{N N}$.

An important role is played by a family of weighted counting functions $\delta_{i}$, which are associated with and defined on the individual $\Gamma_{i}$ (cf. [11, 12, 24, 35, 30]) and are defined for $\gamma \in[1 / 2, \infty)$. Given $\Omega_{i}$ and $\mathbf{x} \in \Gamma_{i, h}, \delta_{i}(\mathbf{x})$ is determined by a sum of contributions from $\Omega_{i}$ and its relevant next neighbors,

$$
\begin{equation*}
\delta_{i}(\mathbf{x})=\sum_{j \in \mathcal{N}_{\mathbf{x}}} \rho_{j}^{\gamma}(\mathbf{x}) / \rho_{i}^{\gamma}(\mathbf{x}), \quad \mathbf{x} \in \Gamma_{i, h} . \tag{15}
\end{equation*}
$$

Here $\mathcal{N}_{\mathbf{x}}, \mathbf{x} \in \Gamma_{h}$, is the set of indices $j$ of the subregions such that $\mathbf{x} \in \Gamma_{j, h}$. The function $\delta_{i}$ is discrete harmonic and thus belongs to $U_{i}$. The pseudoinverses $\delta_{i}^{\dagger} \in U_{i}$ are defined, for $\mathbf{x} \in \Gamma_{i, h}$, by

$$
\begin{equation*}
\delta_{i}^{\dagger}(\mathbf{x})=\delta_{i}^{-1}(\mathbf{x}), \quad \mathbf{x} \in \Gamma_{i, h} \tag{16}
\end{equation*}
$$

We note that these functions provide a partition of unity:

$$
\begin{equation*}
\sum_{i=1}^{N} R_{i}^{T} \delta_{i}^{\dagger}(\mathbf{x}) \equiv 1 \tag{17}
\end{equation*}
$$

In particular, for $u \in U$ we can use the formula

$$
\begin{equation*}
u=\sum_{i=1}^{N} R_{i}^{T} u_{i}, \quad \text { with } u_{i}=\mathcal{H}_{i}\left(\delta_{i}^{\dagger} u\right) \tag{18}
\end{equation*}
$$

Here and from now on, we will tacitly assume that whenever we write $\mathcal{H}_{i}(u v)$ or $\mathcal{H}(u v)$ we first form $I^{k}(u v)$, i.e., map the product of the two functions $u$ and $v$ into the $h p$ finite element space by interpolation, and then extend the result as a discrete harmonic function. If there is no confusion, we will sometime use the notation $u v$ in order to denote $I^{k}(u v)$ or $\mathcal{H}_{i}(u v)$.

A coarse space $U_{0}$ of minimal dimension is defined as

$$
U_{0}=\operatorname{span}\left\{R_{i}^{T} \delta_{i}^{\dagger}\right\} \subset U,
$$

where the span is taken over the floating subdomains. We note that $U_{0}$ consists of piecewise discrete harmonic functions and $R_{0}^{T}$ is the natural injection $U_{0} \subset U$. We consider an exact solver on $U_{0}$

$$
\tilde{s}_{0}(u, v):=a(\mathcal{H} u, \mathcal{H} v)=a(u, v) .
$$

For every substructure $\Omega_{i}$, the local bilinear form is

$$
\tilde{s}_{i}(u, v):=a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u\right), \mathcal{H}_{i}\left(\delta_{i} v\right)\right), \quad u, v \in U_{i} .
$$

For a floating subdomain $\tilde{P}_{i}$ is defined only for those $u \in U$ for which $s(u, v)=$ 0 for all $v=R_{i}^{T} v_{i}$ such that $\mathcal{H}_{i}\left(\delta_{i} v_{i}\right)$ is constant on $\Omega_{i}$. This condition is satisfied if $a\left(u, R_{i}^{T} \delta_{i}^{\dagger}\right)=0$; we note that $R_{i}^{T} \delta_{i}^{\dagger}$ is a basis function for $U_{0}$. For such subdomains, we make the solution $\tilde{P}_{i} u$ of (14) unique by imposing the constraint

$$
\begin{equation*}
\int_{\Omega_{i}} \mathcal{H}_{i}\left(\delta_{i} \tilde{P}_{i} u\right) d \mathbf{x}=0 \tag{19}
\end{equation*}
$$

which just means that we select the solution orthogonal to the null space of the Neumann operator. Thus, Range $\left(\tilde{P}_{i}\right)$ has codimension 1 with respect to the space $U_{i}$.

We can equally well use matrix notations. Let $D_{i}$ be the diagonal matrix with the elements $\delta_{i}^{\dagger}(\mathbf{x})$ corresponding to the point $\mathbf{x} \in \Gamma_{i, h}$. Then

$$
\tilde{s}_{i}(u, v)=u^{T} D_{i}^{-1} S_{i} D_{i}^{-1} v .
$$

We also have,

$$
P_{i}=R_{i}^{T} D_{i} S_{i}^{\dagger} D_{i} R_{i} S
$$

where $S_{i}^{\dagger}$ is a pseudoinverse of $S_{i}$. Analogously for the coarse projection

$$
P_{0}=R_{0}^{T} S_{0}^{-1} R_{0} S
$$

where $S_{0}=R_{0} S R_{0}^{T}$ the restriction of $S$ to $U_{0}$
The main result of this paper is a bound for the condition number of $P_{N N}$. Such bound can be found using the abstract Schwarz theory; see, e.g., [39, Ch. $6]$. We refer to $[24,12,30,39,17]$ for similar proofs.

A uniform bound for the smallest eigenvalue can be found using the decomposition (18) and the fact that $P_{0}$ is an orthogonal projection.

Lemma 6.2 We have

$$
s\left(P_{N N} u, u\right) \geq s(u, u), \quad u \in U
$$

In order to find a bound for the largest eigenvalue, we need a stability property for the local bilinear forms; see, e.g., [39].

Assumption 6.1 We have

$$
s\left(R_{i}^{T} u_{i}, R_{i}^{T} u_{i}\right) \leq \omega \tilde{s}_{i}\left(u_{i}, u_{i}\right), \quad u_{i} \in \operatorname{Range}\left(\tilde{P}_{i}\right), \quad i=1, \cdots, N
$$

with

$$
\omega=C(1-\sigma)^{-6}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}
$$

and $C$ independent of $k, n, \sigma, \gamma$, the coefficients $\rho_{i}$, and the diameters $H_{i}$.
The proof of Assumption 6.1 is given in section 9. Assumption 6.1 and a coloring argument provide a bound for the largest eigenvalue; see, e.g., [30, Sect. 8].
Lemma 6.3 Let Assumption 6.1 be satisfied. Then

$$
s\left(P_{N N} u, u\right) \leq C \omega s(u, u), \quad u \in U
$$

Consequently the condition number of $P_{N N}$ satisfies

$$
\kappa\left(P_{N N}\right) \leq C \omega=C(1-\sigma)^{-6}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}
$$

## 7 Decomposition results

A key ingredient for the proof of Assumption 6.1 and for the analysis of many iterative substructuring methods in three dimensions is a decomposition result for local functions in $U_{i}$ into face and wirebasket components:

$$
\begin{equation*}
u=\sum_{j} u_{F^{i j}}+u_{W^{i}}, \quad u \in U_{i} \tag{20}
\end{equation*}
$$

The face component $u_{F^{i j}}$ vanishes on $\partial \Omega_{i} \backslash F^{i j}$ and is discrete harmonic. It is uniquely determined by the nodal values in $F_{h}^{i j}$. The wirebasket component $u_{W^{i}}$ is also discrete harmonic and vanishes at all points of $\Gamma_{i, h}$ except at those in $W_{h}^{i}$.

We can further decompose a local functions by also defining edge and vertex components:

$$
\begin{equation*}
u=\sum_{j} u_{F^{i j}}+\sum_{j} u_{E^{i j}}+\sum_{j} u_{V^{i j}}, \quad u \in U_{i} \tag{21}
\end{equation*}
$$



Figure 2: Face patch: partition of an edge $E$ that touches $\partial \Omega$ into $E_{1-\sigma}$ and $E_{\sigma}$ (left) and two-dimensional mesh $\mathcal{T}(z)$ for a section corresponding to a constant $z$ (right).
where $u_{E^{i j}}$ is discrete harmonic and vanishes on $\partial \Omega_{i} \backslash E^{i j}$, and $u_{V^{i j}}$ vanishes at all nodes in $\Gamma_{i, h}$ except at the vertex $V^{i j}$. We recall that we exclude geometrical objects on $\partial \Omega$ and that therefore the sums in (20) and (21) are taken over faces, edges, and vertices that do not belong to $\partial \Omega$. Discrete harmonic functions of type $u_{F^{i j}}, u_{E^{i j}}, u_{V^{i j}}$, and $u_{W^{i}}$ are called face, edge, vertex, and wirebasket functions, respectively.

Here and in the following section, we only carry out proofs for the reference cube $\hat{Q}$ : since elements in the macromesh $\mathcal{T}_{m}$ are shape-regular and affinely mapped, the corresponding bounds for a generic substructure $\Omega_{i} \in \mathcal{T}_{m}$, of diameter $H_{i}$, can be obtained by a standard scaling argument and involve the scaled norm (2). We recall that we only need to consider four types of patches: face, edge, corner, and unrefined ones, together with the corresponding triangulations $\mathcal{T}_{f}, \mathcal{T}_{e}, \mathcal{T}_{c}$, and $\hat{Q}$, respectively; cf. Figures 1. We recall that a generic patch is denoted by $\Omega_{i}$ and its triangulation by $\mathcal{T}_{i}$.

### 7.1 Wirebasket components

Given an edge $E=E^{i j} \subset W^{i}$, we define a discrete $L^{2}$ norm on $E$. If $E$ does not touch the boundary $\partial \Omega$, we simply set

$$
\|u\|_{h, E}:=\|u\|_{0, E} .
$$

Let now $E$ be an edge that touches $\partial \Omega$; see Figure 2, left, for an example of a face patch. After a possible translation and rotation, $E$ can always be written as

$$
E=\{(1,1, z) \mid z \in I\} .
$$

Then, the local mesh $\mathcal{T}_{i}$ gives rise to a one-dimensional triangulation on $E, \mathcal{T}_{E}$, which is not quasiuniform and is geometrically refined towards one end point,
say $z=1$. In addition, $E$ can be partitioned as

$$
\bar{E}=\bar{E}_{1-\sigma} \cup \bar{E}_{\sigma}, \quad E_{1-\sigma}=(-1,-1+2(1-\sigma)), \quad E_{\sigma}=(-1+2(1-\sigma), 1)
$$

We note that $E_{1-\sigma}$ consists of exactly one element of length $2(1-\sigma)$ in $\mathcal{T}_{E}$, while the elements on $E_{\sigma}$ are geometrically refined towards $z=1$. We now consider the GLL mesh $\mathcal{T}_{k}\left(\Omega_{i}\right)$ and observe that all the elements that touch the edge $E$ have the same diameters $h_{i, x}$ and $h_{i, y}$, along the two directions perpendicular to $E$; cf. Figure 2. Indeed, $h_{i, x}$ and $h_{i, y}$ are of order $k^{-2}$, for a face patch, of order $k^{-2}(1-\sigma)$, for a corner patch, and of order $k^{-2}$ and $k^{-2}(1-\sigma)$, respectively, for an edge patch. Moreover, thanks to our particular meshes and to the fact that local spaces of the same degree $k$ are employed on each element, we have the following property:

Property 7.1 Let $E$ be an edge parallel to, e.g., z, that is shared by two substructures $\Omega_{i}$ and $\Omega_{j}$. Then, the meshsizes $h_{i, x}$ and $h_{j, x}$, and $h_{i, y}$ and $h_{j, y}$ are comparable. In particular, there exist constants, depending only on the aspect ratios of $\Omega_{i}$ and $\Omega_{j}$, such that

$$
c(1-\sigma) h_{i, x} \leq h_{j, x} \leq C(1-\sigma)^{-1} h_{i, x}
$$

Similar bounds hold for $h_{i, y}$ and $h_{j, y}$.
We define

$$
\|u\|_{h, E}^{2}:=\|u\|_{0, E}^{2}+\|u\|_{h, E_{\sigma}}^{2}=\|u\|_{0, E}^{2}+h_{i, x} h_{i, y}\left\|\partial_{z} u\right\|_{0, E_{\sigma}}^{2} .
$$

We note that in this case the discrete norm is obtained by adding to the $L^{2}$ norm on $E$ a weighted $L^{2}$ norm of $\partial_{z} u$ over a part of $E$ where $\mathcal{T}_{E}$ is not quasiuniform. A discrete wirebasket norm is obtained by summing the contributions over all the edges:

$$
\|u\|_{h, W^{i}}^{2}:=\sum_{E \subset W^{i}}\|u\|_{h, E}^{2}
$$

Lemma 7.1 Let $u_{W^{i}} \in U_{i}$ be discrete harmonic and vanish at all nodal points $\Gamma_{i, h}$ except at those on $W^{i}$. Then there is a constant independent of $u_{W^{i}}, H_{i}$, $\sigma$, and $n$, such that

$$
\left|u_{W^{i}}\right|_{1, \Omega_{i}}^{2} \leq C(1-\sigma)^{-2}\left\|u_{W^{i}}\right\|_{h, W^{i}}^{2}
$$

Proof. The result follows by estimating the energy norm of the zero extension of the boundary values and by noting that the harmonic extension has a smaller energy (cf. Lemma 6.1). More precisely, let $u_{k}$ be the function that vanishes at all nodal points in $\Omega_{i, h} \cup \Gamma_{i, h}$ except at those on $W^{i}$ and $u=u_{h}=I^{h} u_{k}$ the corresponding piecewise trilinear function defined on the GLL mesh $\mathcal{T}_{k}\left(\Omega_{i}\right)$. We will estimate the energy of $u_{h}$ on each element $K \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ that touch an edge
$E \subset W^{i}$. Without loss of generality, we assume that $E$ is parallel to the $z$ axis. We only consider the worst possible case, i.e., that of a face patch and refer to Figure 2.

Let us first suppose that $E$ does not touch $\partial \Omega$. For a face patch, $K$ has dimensions $h_{x}, h_{y}$, and $h_{z}$ of order

$$
\begin{aligned}
& k^{-2} \times k^{-2}(1-\sigma) \times k^{-2}, \quad \text { or } \\
& k^{-2} \times k^{-2}(1-\sigma) \times k^{-1},
\end{aligned}
$$

and thus

$$
\begin{align*}
& c(1-\sigma) h_{x} \leq h_{y} \leq C h_{x}  \tag{22}\\
& h_{x} \leq C h_{z}
\end{align*}
$$

see Figure 2. If $a$ and $b$ are the vertices of $K$ that lie on $E$, Lemma 5.3 yields

$$
\left\|\partial_{x} u\right\|_{0, K}^{2} \leq C\left(h_{y} h_{z} / h_{x}\right)\left(u(a)^{2}+u(b)^{2}\right) \leq C \int_{a}^{b} u^{2} d z
$$

where for the last inequality we have used (22) and standard properties of linear functions. In a similar way, we find

$$
\left\|\partial_{y} u\right\|_{0, K}^{2} \leq C(1-\sigma)^{-1} \int_{a}^{b} u^{2} d z, \quad\left\|\partial_{z} u\right\|_{0, K}^{2} \leq C \int_{a}^{b} u^{2} d z
$$

Let now $E$ be an edge that touches $\partial \Omega$ and $K \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ be an element that shares an edge with $E_{1-\sigma}$. For a face patch, $K$ has dimensions of the order

$$
\begin{aligned}
& k^{-2} \times k^{-2} \times k^{-2}(1-\sigma), \quad \text { or } \\
& k^{-2} \times k^{-2} \times k^{-1}(1-\sigma),
\end{aligned}
$$

and thus

$$
\begin{align*}
& c h_{x} \leq h_{y} \leq C h_{x} \\
& h_{x} \leq C(1-\sigma)^{-1} h_{z} \tag{23}
\end{align*}
$$

see Figure 2, left. As before, Lemma 5.3 yields

$$
\left\|\partial_{x} u\right\|_{0, K}^{2} \leq C \int_{a}^{b} u^{2} d z, \quad\left\|\partial_{y} u\right\|_{0, K}^{2} \leq C \int_{a}^{b} u^{2} d z, \quad\left\|\partial_{z} u\right\|_{0, K}^{2} \leq C(1-\sigma)^{-2} \int_{a}^{b} u^{2} d z
$$

We are now left with the case of an element $K \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ that shares an edge with $E_{\sigma}$. We note that the first of (23) remains valid in this case. We then have

$$
\left\|\partial_{x} u\right\|_{0, K}^{2} \leq C \int_{a}^{b} u^{2} d z, \quad\left\|\partial_{y} u\right\|_{0, K}^{2} \leq C \int_{a}^{b} u^{2} d z
$$

For $\partial_{z} u$, we trivially have

$$
\left\|\partial_{z} u\right\|_{0, K}^{2} \leq C\left(h_{x} h_{y} / h_{z}\right)(u(a)-u(b))^{2} \leq C h_{x} h_{y} \int_{a}^{b}\left(\partial_{z} u\right)^{2} d z
$$

The proof is concluded by summing over the elements $K \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ and using Lemma 5.1.

We now have a bound for the wirebasket component.
Theorem 7.1 Let $u \in U_{i}$ and $u_{W^{i}}$ be its wirebasket component. Then there is a constant independent of $u, H_{i}, \sigma$, and $n$, such that

$$
\left|u_{W^{i}}\right|_{1, \Omega_{i}}^{2} \leq C(1-\sigma)^{-2}\|u\|_{h, W^{i}}^{2}
$$

A complementary result is given by the trace estimates in Lemma 7.2. We first introduce some additional notations. Let $E$ be an edge of a substructure $\Omega_{i}$. Without loss of generality, we assume that $\Omega_{i}$ coincides with the reference cube $\hat{Q}$ and that $E=\{(1,1, z) \mid z \in I\}$. The intersection between the plane corresponding to a constant $z \in I$ and $\hat{Q}$ is the unit square $\hat{S}=(-1,1)^{2}$, and the local mesh $\mathcal{T}_{i}$ gives rise to a two-dimensional mesh $\mathcal{T}(z)$ on $\hat{S}$ which is either a two-dimensional edge or corner patch, or it consists of a single element $\hat{S}$; see Figure 2, right. Let $V=(1,1)$ be the intersection between $E$ and the closure of $\hat{S}$. If $K_{V} \in \mathcal{T}(z)$ is the two-dimensional element that contains $V$, we note that, since $E$ does not belong to $\partial \Omega, K_{V}$ has dimensions in $\{2,2(1-\sigma)\}$, and thus independent of the level of refinement $n$. For a fixed $(x, y) \in \bar{K}_{V}$, we finally define the edge $E(x, y)=\{(x, y, z) \mid z \in I\}$.

Lemma 7.2 Let $u_{k} \in X_{i}$ and $E$ and edge of $\Omega_{i}$. Then there is a constant independent of $u_{k}, H_{i}, \sigma$, and $n$, such that

$$
\begin{aligned}
& \left\|u_{k}\right\|_{0, E}^{2} \leq C(1-\sigma)^{-2}(1+\log k)\left\|u_{k}\right\|_{1, \Omega_{i}}^{2} \\
& \left\|u_{k}\right\|_{h, E}^{2} \leq C(1-\sigma)^{-2}(1+\log k)\left\|u_{k}\right\|_{1, \Omega_{i}}^{2}
\end{aligned}
$$

Proof. As before, it is enough to find bounds for $u=I^{h} u_{k}$. Without loss of generality, we assume $E=\{(1,1, z) \mid z \in I\}$. We consider the two-dimensional mesh $\mathcal{T}(z)$ on the intersection between the plane corresponding to a constant $z$ and the substructure; cf. Figure 2, right. Since geometric refinement on $\mathcal{T}(z)$ takes place far from the vertex $(1,1)$, we can apply the two-dimensional result in [40, Lem. 7.6] and write

$$
|u(1,1, z)|^{2} \leq C(1-\sigma)^{-2}(1+\log k)\|u(\cdot, \cdot, z)\|_{1, \hat{S}}^{2}, \quad z \in(-1,1)
$$

with a constant that is independent of $n, \sigma$, and $z$. Integrating over $z$ then gives

$$
\|u\|_{0, E}^{2} \leq C(1-\sigma)^{-2}(1+\log k)\|u\|_{1, \Omega_{i}}^{2},
$$

which proves the first inequality and the second one for edges that do not touch $\partial \Omega$.

We now bound $\|u\|_{h, E_{\sigma}}$ for an edge that touches the boundary $\partial \Omega$. We consider the one-dimensional GLL meshes for each one of the elements in $\mathcal{T}_{E}$ and estimate the single contributions from the elements of these meshes. Let $e$
be one of these elements of length $h_{z}$ and end points $a$ and $b$. The edge $e$ belongs to a parallelepiped $K_{e} \in \mathcal{T}_{k}\left(\Omega_{i}\right)$. We note that $K_{e}$ has dimensions $h_{x}=h_{i, x}$, $h_{y}=h_{i, y}$, and $h_{z}$. Since $u$ is linear on $e$ and trilinear on $K_{e}$, we have

$$
h_{x} h_{y} \int_{e} \partial_{z} u^{2} d z \leq C \frac{h_{x} h_{y}}{h_{z}}(u(a)-u(b))^{2} \leq C\left\|\partial_{z} u\right\|_{0, K_{e}}^{2},
$$

where, for the last inequality, we have used Lemma 5.3. Summing over the edges $e$ in $E_{\sigma}$ yields

$$
\|u\|_{h, E_{\sigma}}^{2} \leq C\left\|\partial_{z} u\right\|_{0, \Omega_{i}}^{2}
$$

which, combined with the first inequality, proves the second bound.
The next lemma can be proved using the two-dimensional bound in [40, Lem. 7.6] and similar arguments as before. We note that it is only valid for edges $E(x, y)$ that are not too far from $E$ and thus not too close to the part of $\Omega_{i}$ where anisotropic refinement takes place.

Lemma 7.3 Let $E$ be an edge of a substructure $\Omega_{i}$ which is parallel, say, to $z$ and intersects the plane corresponding to a constant $z$ in $V$. Let in addition $K_{V}$ be the element in the two-dimensional mesh $\mathcal{T}(z)$ that contains $V$. Then, for every $(x, y) \in \bar{K}_{V}$ and $u_{k} \in X_{i}$,

$$
\begin{equation*}
\left\|u_{k}\right\|_{0, E(x, y)}^{2} \leq C(1-\sigma)^{-2}(1+\log k)\left\|u_{k}\right\|_{1, \Omega_{i}}^{2} \tag{24}
\end{equation*}
$$

where $C$ is independent of $u_{k}, \sigma, n, k$, and $(x, y)$, but depends only on the aspect ratio of $\Omega_{i}$.

Proof. The proof can be carried out as in the previous lemma by using the two-dimensional result in [40, Lem. 7.6]. Indeed, since the point $(x, y)$ belongs to $\bar{K}_{V}$ and is thus far from the region where anisotropic refinement takes place, we have

$$
|u(x, y, z)|^{2} \leq C(1-\sigma)^{-2}(1+\log k)\|u(\cdot, \cdot, z)\|_{1, \hat{S}}^{2}, \quad z \in(-1,1)
$$

Integration along $z$ concludes the proof. D
We end this section with a stability result for vertex and edge components. It is a direct consequence of (9) and of the fact that for a vertex function the modified norm $\|\cdot\|_{h, E}$ coincides with $\|\cdot\|_{0, E}$.

Lemma 7.4 Let $E$ be an edge of a substructure $\Omega_{i}$ and $V$ one of its end points. Then, for every $u \in X_{i}$,

$$
\begin{equation*}
\left\|u_{V}\right\|_{h, W^{i}}^{2} \leq C\|u\|_{h, W^{i}}^{2}, \quad\left\|u_{E}\right\|_{h, W^{i}}^{2} \leq C\|u\|_{h, W^{i}}^{2} \tag{25}
\end{equation*}
$$

where $C$ is independent of $u, \sigma, n, k$.


Figure 3: Edge patch on the reference cube $(-1,1)^{3}$ employed in the proofs of Lemmas 7.5 and 7.6

### 7.2 Face components

We next consider the face contributions of the decomposition (20). Bounds for face contributions on the unrefined patch follow from standard results for spectral elements. For face, edge, and corner patches, we employ cut-off functions $\theta_{F}$ for each face and Lemma 5.2. We note that we need to consider one possible case for faces of the corner patch, and two for the edge and face patches; cf. Figure 1. In this section we only consider the case of an edge patch $\Omega_{i}$ in full detail, with the edge $(1, y,-1), y \in I$, and the two adjacent faces in common with $\partial \Omega$; see Figure 3. The other patches can be dealt with in a similar way.

As shown if Figure 3 for the reference cube, the edges that do not lie on $\partial \Omega$ are denoted by $E^{l}, l=1, \ldots, 5$, with $E^{5}$ the edge that does not touch the boundary $\partial \Omega$. An edge patch is further partitioned into three regions. The first step of geometric refinement partitions $\hat{Q}$ into four parallelepipeds with dimensions in $\{2,2(1-\sigma), 2 \sigma\}$. Let $K_{\Omega}$ be the one that contains the boundary edge and $K_{\text {int }}$ the one that does not touch $\partial \Omega$ and contains the inner edge $E^{5}$. The two remaining parallelepipeds are denoted by $K^{12}$ and $K^{34}$ and they touch the edges $E^{1}$ and $E^{2}$, and $E^{3}$ and $E^{4}$, respectively. The region $K_{\text {edge }}$ is the union of $K^{12}$ and $K^{34}$; cf. Figure 3.

The proof of the following lemma is a modification of those of [10, Lem. 3.3.6] and [40, Lem. 7.7].

Lemma 7.5 Given a face $F^{j}$ of $\Omega_{i}$ that does not lie on $\partial \Omega$, there exists a continuous function $\theta_{F^{j}}$, defined on $\overline{\Omega_{i}}$, that is equal to one at the nodal points
of $F_{h}^{j}$ and zero on $\Gamma_{i, h} \backslash F_{h}^{j}$, such that

$$
\begin{align*}
& \sum_{F^{j} \subset \Gamma_{i}} \theta_{F^{j}}(\mathbf{x})=1, \quad \mathbf{x} \in\left(\Omega_{i, h} \cup \Gamma_{i, h}\right) \backslash W_{h}^{i}, \\
& 0 \leq \theta_{F^{j}} \leq 1,  \tag{26}\\
& \left|\nabla \theta_{F^{j}}\right| \leq C / r, \quad \text { in } \Omega_{i} \backslash K_{\Omega} \\
& \left|\nabla \theta_{F^{j}}\right| \leq C / H_{i}, \quad \text { in } K_{\Omega},
\end{align*}
$$

where $r=r(\mathbf{x})$ is the distance to the closest edge of $\Omega_{i}$ that does not lie on $\partial \Omega$.
Proof. We only need to construct four functions and we will do that by constructing them in the three regions $K_{\text {int }}, K_{\text {edge }}$, and $K_{\Omega}$ separately.

We start with the inner region $K_{\text {int }}$ and employ a similar construction as in [10, Lem. 3.3.6]:
We further partition $\Omega_{i}$ into eight parallelepipeds by bisecting $\left\{K_{i n t}, K^{12}, K^{34}, K_{\Omega}\right\}$ with the plane $y=0$; see Figure 3, left. Let the center $C$ be the common vertex to these parallelepipeds and $\left\{C^{j}, j=1, \ldots, 6\right\}$, be their vertices that belong to the six faces of $\Omega_{i}$; see Figure 3, right. By connecting the center $C$ with the centers $C^{j}$ and with the eight vertices of $\Omega_{i}$, and, for each face, by connecting the point $C^{j}$ with the four vertices of this face, we can partition $\Omega_{i}$ into twenty-four tetrahedra; see Figure 3, right. By intersecting them with $K_{i n t}$, we obtain a partition of $K_{\text {int }}$ into eight tetrahedra. We first define a function $\vartheta_{F^{j}}$ associated to the face $F^{j}$, defined to be $1 / 4$ at the center $C$ and $\vartheta_{F^{j}}\left(C^{l}\right)=\delta_{j l}$ at the centers of the faces. On the segments $C C^{l}$, these functions are obtained by linear interpolation of the values at $C$ and $C^{l}$; see Figure 3, right. The values inside each subtetrahedron formed by the segment $C C^{l}$ and one edge of $F^{l}$ are defined to be constant on the intersection of any plane through that edge, and are given by the value on the segment $C C^{l}$. We note that this procedure determines $\vartheta_{F^{j}}$ at all points in $\Omega_{i}$ except on the wirebasket $W^{i}$.

We next consider the GLL triangulation $\mathcal{T}_{k}\left(\Omega_{i}\right)$ and interpolate $\vartheta_{F^{j}}$ at the GLL nodes in $\bar{K}_{\text {int }} \backslash W^{i}$ :

$$
\theta_{F^{j}}(\mathbf{x})=\left(I^{h} \vartheta_{F^{j}}\right)(\mathbf{x}), \quad \mathbf{x} \in \bar{K}_{i n t} \backslash W^{i} .
$$

The function $\theta_{F^{j}}$ is set to zero on the nodes in $W_{h}^{i}$. The functions $\theta_{F^{j}}$ are non negative and bounded by one: this proves the second of (26) for points in $K_{\text {int }}$. By construction, also the first of (26) holds for every node in $\bar{K}_{i n t} \backslash W^{i}$. The third of (26) can be proven by proceeding in the same way as for [10, Lem. 3.3.6].

We next construct the functions $\theta_{F^{j}}$ in $K_{\text {edge }}$ :
We start with $K^{12}$. We take the values on the common face $\overline{K^{12}} \cap \bar{K}_{\text {int }}$ and we extend them as constants into $\overline{K^{12}}$ along the segments parallel to $E^{1}$ and $E^{2}$; see Figure 3, left. The inequalities in (26) remain valid. We note that the
function obtained is independent of $x$ in $K^{12}$. A similar construction is carried out in $K^{34}$.

Finally we construct $\theta_{F^{j}}$ in $K_{\Omega}$ :
We note that $K_{\Omega}$ is divided into two parallelepipeds and that on their internal faces the function $\theta_{F^{j}}$ has already been defined. In addition, $\theta_{F^{j}}$ is bilinear on these faces. It is then enough to assign the value $1 / 4$ at the end points and mid point of the boundary edge and interpolate these values in $K_{\Omega}$ in order to obtain a piecewise trilinear function. The first, second and fourth of (26) follow from standard properties of trilinear functions.

By examining the proof of the previous lemma, we see that, for an edge $E$ that touches $\partial \Omega$, the value of the functions $\theta_{F^{j}}$ is independent of the coordinate along the direction of $E$ in all the elements of the GLL meshes that touch $E_{\sigma}$; cf. Figure 3, left.

Property 7.2 Let $F$ be a face of $\Omega_{i}$ and $E$ be an edge, parallel to, say z, that touches $\partial \Omega$. In any element $K_{E} \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ that shares an edge with $E_{\sigma}$ the function $\theta_{F}$ is independent of $z$.

We are now able to bound the face components in the decomposition (20).
Lemma 7.6 Let $\theta_{F^{j}}$ be the functions in Lemma 7.5, where $F^{j}$ is a face of the substructure $\Omega_{i}$. Then, for every $\mathbf{x} \in \Omega_{i, h} \cup \Gamma_{i, h}$ that is not on the wirebasket of $\Omega_{i}$,

$$
\sum_{j} I^{k}\left(\theta_{F^{j}} u\right)(\mathbf{x})=\sum_{j} I^{h}\left(\theta_{F^{j}} u\right)(\mathbf{x})=u(\mathbf{x}), \quad u \in X_{i}
$$

and

$$
\left|I^{k}\left(\theta_{F^{j}} u\right)\right|_{1, \Omega_{i}}^{2} \leq C(1-\sigma)^{-4}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}\|u\|_{1, \Omega_{i}}^{2}
$$

Proof. We only consider the case of an edge patch $\Omega_{i}$ in full detail; see Figure 3. The proof is similar to that in [40, Lem. 7.8] and [10, Lem. 3.3.7] but particular care is required close to the edges that touch $\partial \Omega$. Indeed, thanks to Lemma 5.1, it is enough to find a bound for the piecewise trilinear function $I^{h}\left(\theta_{F^{j}} u\right)$.

The first equality follows directly from the first of (26). For the second inequality, we consider an element $K$, of dimensions $h_{x}, h_{y}$, and $h_{z}$, in the GLL mesh $\mathcal{T}_{k}\left(\Omega_{i}\right)$. We consider three cases (as opposed to [10, Lem. 3.3.7] where only two cases are considered): $K$ may belong to the region $K_{\Omega}$ containing the boundary edge, touch the wirebasket, or may not touch it; see Figure 3.

Case 1. We start with an element that touches an edge $E$ and does not belong to $K_{\Omega}$. We can proceed as in [10, Lem. 3.3.7] if $E$ does not touch $\partial \Omega$ $\left(E=E^{5}\right)$ or, in case it does $\left(E=E^{l}, l=1, \ldots, 4\right)$, if $K$ does not touch $E_{\sigma}$. We only consider the case of $E=E^{3}$ in full detail; cf. Figure 3, left. The nodal values of $I^{h}\left(\theta_{F^{j}} u\right)$ on $K$ are $0,0,0,0, u(a), u(b), \theta_{F^{j}}(c) u(c)$, and $\theta_{F^{j}}(d) u(d)$,
with $a$ and $b$ vertices on a face and $c$ and $d$ vertices inside $\Omega_{i}$. It is immediate to see that

$$
\begin{align*}
& c(1-\sigma) h_{x} \leq h_{y} \leq C(1-\sigma)^{-1} h_{x} \\
& h_{x} \leq C(1-\sigma)^{-1} h_{z} \tag{27}
\end{align*}
$$

Using Lemma 5.3 and (27), we can easily find

$$
\begin{aligned}
\mid I^{h}\left(\left.\theta_{F^{j}} u\right|_{1, K} ^{2}\right. & \leq C(1-\sigma)^{-2} h_{z}\left(u(a)^{2}+u(b)^{2}+u(c)^{2}+u(d)^{2}\right) \\
& \leq C(1-\sigma)^{-2}\left(\int_{a}^{b} u^{2} d z+\int_{c}^{d} u^{2} d z\right)
\end{aligned}
$$

where we have also used the fact that $\theta_{F^{j}}$ has values between zero and one. Summing over the element $K$ and using in Lemma 7.3 for segments that are parallel to $E$ give

$$
\sum_{K}\left|I^{h}\left(\theta_{F^{j}} u\right)\right|_{1, K}^{2} \leq C(1-\sigma)^{-4}(1+\log k)\|u\|_{1, \Omega_{i}}^{2}
$$

where the sum is taken over the elements in $\mathcal{T}_{k}\left(\Omega_{i}\right)$ that touch an edge $E$, such that $E$ does not touch $\partial \Omega$ or, if it does, $K$ does not touch $E_{\sigma}$.

We next consider the case where $K$ shares an edge with $E_{\sigma}$. The terms involving the $x$ and $y$ derivatives can be bounded as before: indeed, the first of (27) still holds in this case. However, the second of (27), needed to bound the $z$ derivative, does not hold. Using Lemma 5.3 we find
$\left\|\partial_{z} I^{h}\left(\theta_{F^{j}} u\right)\right\|_{0, K}^{2} \leq C\left(h_{x} h_{y} / h_{z}\right)\left((u(a)-u(b))^{2}+\left(\theta_{F^{j}}(c) u(d)-\theta_{F^{j}}(d) u(d)\right)^{2}\right)$.
Property 7.1 ensures that $\theta_{F^{j}}(c)=\theta_{F^{j}}(d)$ and thus

$$
\left\|\partial_{z} I^{h}\left(\theta_{F^{j}} u\right)\right\|_{0, K}^{2} \leq C\left\|\partial_{z}\left(\theta_{F^{j}} u\right)\right\|_{0, K}^{2}
$$

Summing over the elements $K$ that touch $E_{\sigma}$ gives

$$
\sum_{K} \| \partial_{z}\left(I^{h}\left(\theta_{F^{j}} u\right)\left\|_{0, K}^{2} \leq C\right\| \partial_{z}\left(\theta_{F^{j}} u\right) \|_{0, \Omega_{i}}^{2}\right.
$$

and thus

$$
\begin{equation*}
\sum_{\bar{K} \cap W^{i} \neq \emptyset}\left|I^{h}\left(\theta_{F^{j}} u\right)\right|_{1, K}^{2} \leq C(1-\sigma)^{-4}(1+\log k)\|u\|_{1, \Omega_{i}}^{2} . \tag{28}
\end{equation*}
$$

Case 2. We now consider an element $K \in \mathcal{T}_{k}\left(\Omega_{i}\right)$ that does not touch the wirebasket and does not belong to $K_{\Omega}$. The proof for this case is similar to that of [10, Lem. 3.3.7]. Using Lemma 5.2 and the second of (26), we have

$$
\sum_{\substack{K \subset \Omega_{i} \backslash K_{\Omega} \\ \bar{K} \cap W^{i}=\emptyset}}\left|I^{h}\left(\theta_{F^{j}} u\right)\right|_{1, K}^{2} \leq C \sum_{K}\left(|u|_{1, K}^{2}+r_{K}^{-2}\|u\|_{0, K}^{2}\right),
$$

where $r_{K}$ is the distance of the baricenter of $K$ from the wirebasket. We have

$$
\begin{aligned}
& \sum_{K} r_{K}^{-2}\|u\|_{0, K}^{2} \leq C \int_{K_{\text {int }} \cup K^{12} \cup K^{34}} r^{-2} u^{2} d \mathbf{x} \\
& \leq C \int_{K_{\text {int }}} r_{5}^{-2} u^{2} d \mathbf{x}+C \sum_{j=1}^{2} \int_{K^{12} \cup K_{\text {int }}} r_{j}^{-2} u^{2} d \mathbf{x}+C \sum_{j=3}^{4} \int_{K^{34} \cup K_{\text {int }}} r_{j}^{-2} u^{2} d \mathbf{x}
\end{aligned}
$$

where $r_{j}$ denotes the distance of a point from the edge $E^{j}$ and the region consisting of the elements in the GLL mesh $\mathcal{T}_{k}\left(\Omega_{i}\right)$ that touch the wirebasket is assumed to be excluded from the domains of integration; cf. Figure 3, left. Each of the integrals on the right, associated to an edge $E=E^{j}$, can be estimated using cylindrical coordinates with the $\zeta$ axis coinciding with $E^{j}$ and the radial direction $r_{j}$ normal to $E^{j}$. We only consider $E^{5}$ in detail; cf. Figure 3. The other integrals can be estimated in the same way. If the point $V$ is the intersection between $E^{5}$ and the section corresponding to a fixed $\zeta$ and $K_{V}$ is the element of the two-dimensional mesh $\mathcal{T}(\zeta)$ that contains $V$, we can write

$$
\begin{aligned}
\int_{K_{\text {int }}} r_{5}^{-2} u^{2} d \mathbf{x} & \leq C \int_{K_{V}} r_{5}^{-2} d x d y \int_{-1}^{1} u^{2} d \zeta \\
& \leq C(1-\sigma)^{-2}(1+\log k)\|u\|_{1, \Omega_{i}}^{2} \int_{K_{V}} r_{5}^{-2} d x d y
\end{aligned}
$$

where we have used Lemma 7.3 for the last inequality; cf. Figure 2, right. The last integral can be estimated by

$$
\int_{K_{V}} r_{5}^{-2} d x d y \leq C \int_{k^{-2}(1-\sigma)}^{2} r_{5}^{-1} d r_{5} \int_{0}^{2 \pi} d \phi \leq C\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)
$$

Considering similar contributions for the other edges, we then find

$$
\begin{equation*}
\sum_{\substack{K \subset \Omega_{i} \backslash K_{\Omega} \\ \bar{K} \cap W^{i}=\emptyset}}\left|I^{h}\left(\theta_{F^{j}} u\right)\right|_{1, K}^{2} \leq C|u|_{1, \Omega_{i}}^{2}+C(1-\sigma)^{-2}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}\|u\|_{1, \Omega_{i}}^{2} \tag{29}
\end{equation*}
$$

Case 3. We are now left with the case $K \subset K_{\Omega}$. Since, in this case, $\left|\nabla \theta_{F^{j}}\right|$ is bounded by a constant, Lemma 5.2 ensures

$$
\sum_{K \subset K_{\Omega}}\left|I^{h}\left(\theta_{F^{j}} u\right)\right|_{1, K}^{2} \leq C\|u\|_{1, \Omega_{i}}^{2}
$$

The proof is concluded by combining this inequality with (28) and (29), and applying Lemma 5.1.


Figure 4: The cross sections of an edge and a face patch, or a corner and an edge patch, with a common face $F$.

## 8 Comparison results

In the analysis of many iterative substructuring methods, it is necessary to compare certain norms of discrete harmonic functions on different substructures that have the same trace on a common face, edge, or vertex.

As already pointed out in [40], if the local meshes are shape-regular and quasi-uniform, the comparison for functions on adjacent substructures that have the same value on a common face, can be made using a trace theorem (which is valid for general functions in $H^{1}$ ) and a stable extension. However, the existence of stable extensions for meshes that are not quasi-uniform or shape-regular is far from trivial. For this reason, here we will adopt the same strategy as in [40], since the meshes considered are highly anisotropic but of a particular type.

We note that we only need to consider three cases: that of a face shared by an unrefined and a face patch, by a face and an edge patch, and by an edge and a corner patch. We only consider the last two case in full detail, since the former can be treated in exactly the same way. We consider the two substructures $\Omega_{i}$ and $\Omega_{j}$ in Figure 4 , which share the face $F$. Since we proceed in exactly the same way as in [40, Sect. 7.3], we do not present any proof here. We first consider $\Omega_{i}$ and suppose that it coincides with the reference cube $\hat{Q}$. The face $F$ corresponds to $x=1$. Let $\Omega_{F}$ be the layer of points in $\Omega_{i}$ within a distance $2(1-\sigma)$ from $F$.

The following lemma can be proven in the same way as [40, Lem. 7.9].
Lemma 8.1 Let $u_{F} \in U_{i}$ be a face function on $\Omega_{i}$, i.e, a discrete harmonic function that vanishes on $\partial \Omega_{i} \backslash F$, and $\tilde{u}_{F} \in X_{i}$, such that

1. $\tilde{u}_{F}$ is equal to $u_{F}$ on $F$ and vanishes on $\partial \Omega_{F} \backslash F$;
2. $\tilde{u}_{F}$ is discrete harmonic in $\Omega_{F}$;
3. $\tilde{u}_{F}$ vanishes in $\Omega_{i} \backslash \Omega_{F}$.

Then

$$
\left|u_{F}\right|_{1, \Omega_{i}}^{2} \leq\left|\tilde{u}_{F}\right|_{1, \Omega_{i}}^{2} \leq\left\|\nabla \theta_{\sigma, F}\right\|_{\infty}^{2}\left|u_{F}\right|_{1, \Omega_{i}}^{2}
$$

where $\theta_{\sigma, F} \in W^{1, \infty}\left(\Omega_{i}\right)$ is any function that is equal to one on $F$, vanishes in $\Omega_{i} \backslash \Omega_{F}$, and has values in $(0,1)$ in the rest of $\Omega_{i}$. In particular we can find a function such that

$$
\left\|\nabla \theta_{\sigma, F}\right\|_{\infty} \leq C(1-\sigma)^{-1}
$$

The comparison result for face functions can be then found by noting that we can map $\Omega_{j}$ and its mesh into $\Omega_{F}$ and the corresponding local mesh, by a simple dilation in the horizontal direction.

Theorem 8.1 Let $F$ be a face that is common to $\Omega_{i}$ and $\Omega_{j}$ and $u_{F} \in U$ be a piecewise discrete harmonic function that is identically zero at all nodal points in $\Gamma_{h} \backslash F_{h}$. Then,

$$
c(1-\sigma)\left|u_{F}\right|_{1, \Omega_{i}}^{2} \leq\left|u_{F}\right|_{1, \Omega_{j}}^{2} \leq C(1-\sigma)^{-1}\left|u_{F}\right|_{1, \Omega_{i}}^{2}
$$

For vertex and edge functions the following lemma is sufficient for our analysis.

Lemma 8.2 Let $\Omega_{i}$ and $\Omega_{j}$ be two substructures and $u \in X$. If $V=V^{i}=V^{j}$ is a common vertex, then the vertex components of $u$ satisfy

$$
\left\|u_{V^{j}}\right\|_{h, W^{j}}^{2} \leq C(1-\sigma)^{-1}\left\|u_{V^{i}}\right\|_{h, W^{i}}^{2}
$$

If $E=E^{i}=E^{j}$ is a common edge, then the edge components of $u$ satisfy

$$
\left\|u_{E^{j}}\right\|_{h, W^{j}}^{2} \leq C(1-\sigma)^{-2}\left\|u_{E^{i}}\right\|_{h, W^{i}}^{2}
$$

Proof. For the first inequality, we note that the modified norms $\|\cdot\|_{h, W^{i}}$ and $\|\cdot\|_{h, W^{j}}$ coincide with the $L^{2}$ norms, since a vertex function vanishes at all nodal points in $\Gamma_{h}$ except at that vertex and we only consider internal vertices. It is enough to compare a contribution from an edge $E^{j}$ of $\Omega_{j}$ with that of an edge $E^{i}$ of $\Omega_{i}$. The worst possible case occurs when $E^{j}$ does not touch $\partial \Omega$ but $E^{i}$ does; cf. Figure 4. Let $\phi(\hat{z})$ be the function in $\mathbb{Q}_{k}(I)$ that vanishes at all the GLL nodes in $I$, except at -1 where it is equal to $u(V)$. Using the change of variables $z=(1-\sigma)(\hat{z}+1)-1$ and the fact that $u_{V^{i}}$ vanishes in $E_{\sigma}^{i}$, we have

$$
\begin{aligned}
\int_{E^{j}} u_{V^{j}}(\hat{z})^{2} d \hat{z} & =\int_{-1}^{1} \phi(\hat{z})^{2} d \hat{z}=(1-\sigma)^{-1} \int_{-1}^{-1+2(1-\sigma)} \phi(z)^{2} d z \\
& =(1-\sigma)^{-1} \int_{E_{1-\sigma}^{i}} u_{V^{i}}(z)^{2} d z=(1-\sigma)^{-1} \int_{E^{i}} u_{V^{i}}(z)^{2} d z
\end{aligned}
$$

For the second inequality, it is enough to use the definition of the modified norms $\|\cdot\|_{h, W^{i}}$ and $\|\cdot\|_{h, W^{j}}$ and Property 7.1 ■

## 9 Proof of Assumption 6.1

We are now ready to give an upper bound for $\omega$ in Assumption 6.1. Our proof is similar to that in [31, Lem. 9.1]. We note that if $u_{i} \in U_{i}$, its extension $u=R_{i}^{T} u_{i}$ vanishes on $\Gamma_{h}$ except at the nodal points in $\Gamma_{i, h}$ and its support is thus contained in the union of $\Omega_{i}$ and its neighboring substructures. In order to estimate $\omega$ we thus have to estimate the energy of $u$ in these substructures in terms of the energy of $\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)$ in $\Omega_{i}$ alone.

We first note that, by simple calculation, we have

$$
\begin{equation*}
\rho_{j}\left(\delta_{i}^{\dagger}(\mathbf{x})\right)^{2}=\rho_{j} \delta_{i}(\mathbf{x})^{-2} \leq \min \left\{\rho_{i}, \rho_{j}\right\}, \quad \mathbf{x} \in \Gamma_{i, h}, \quad j \in \mathcal{N}_{\mathbf{x}} \tag{30}
\end{equation*}
$$

Let $u_{i} \in \operatorname{Range}\left(\tilde{P}_{i}\right)$. We start with a substructure $\Omega_{j}$ that only has a vertex $V=V^{i}=V^{j}$ in common with $\Omega_{i}$. We note that, according to the decomposition (21), $u$ has only a wirebasket component $u_{V^{j}}=u$ on $\Omega_{j}$, which vanishes at all nodes in $\Gamma_{j, h}$ except at $V$. Using Lemma 7.1, we find

$$
a_{j}(u, u)=\rho_{j}\left|u_{V^{j}}\right|_{1, \Omega_{j}}^{2} \leq C \rho_{j}(1-\sigma)^{-2}\left\|u_{V^{j}}\right\|_{h, W^{j}}^{2}=C \rho_{j} \delta_{i, V}^{-2}(1-\sigma)^{-2}\left\|\delta_{i} u_{V^{j}}\right\|_{h, W^{j}}^{2}
$$

where $\delta_{i, V}=\delta_{i}(V)$. We next note that, thanks to Lemma 8.2, the norm $\|\cdot\|_{h, W^{j}}$ associated to $\Omega_{j}$ can be bounded by $\|\cdot\|_{h, W^{i}}$. In addition, we can apply Lemmas 7.4 and 7.2 and find

$$
\begin{aligned}
& \rho_{i}\left\|\delta_{i} u_{V^{j}}\right\|_{h, W^{j}}^{2} \leq C(1-\sigma)^{-1} \rho_{i}\left\|\left(\delta_{i} u_{i}\right)_{V^{i}}\right\|_{h, W^{i}}^{2} \leq C(1-\sigma)^{-1} \rho_{i}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{h, W^{i}}^{2} \\
& \quad \leq C(1-\sigma)^{-3}(1+\log k) \rho_{i}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{1, \Omega_{i}}^{2} \\
& \quad=C(1-\sigma)^{-3}(1+\log k)\left(a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u_{i}\right), \mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right)+\rho_{i} H_{i}^{-2}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{0, \Omega_{i}}^{2}\right) .
\end{aligned}
$$

The $L^{2}$ component in the last term can be bounded by the local bilinear form $a_{i}(\cdot, \cdot)$, thanks to a Poincaré inequality for floating subdomains (cf. (19)), or thanks to a Friedrichs inequality for substructures that touch $\partial \Omega$. Combining these two estimates and using (30), we find

$$
\begin{equation*}
a_{j}(u, u)=a_{j}\left(u_{V^{j}}, u_{V^{j}}\right) \leq C(1-\sigma)^{-5}(1+\log k) a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u_{i}\right), \mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right) \tag{31}
\end{equation*}
$$

We next consider a substructure $\Omega_{j}$ that only has an edge $E=E^{i}=E^{j}$ in common with $\Omega_{i}$, with vertices $V^{j 1}=V^{i 1}$ and $V^{j 2}=V^{i 2}$. We note that, according to the decompositions (20) and (21), $u$ has only a wirebasket component on $\Omega_{j}$,

$$
u=u_{W^{j}}=u_{V^{j 1}}+u_{V^{j 2}}+u_{E^{j}}
$$

which vanishes at all nodes in $\Gamma_{j, h}$ except at those on the closure $\overline{E^{j}}$. We then have

$$
a_{j}(u, u) \leq 3 a_{j}\left(u_{V^{j 1}}, u_{V^{j 1}}\right)+3 a_{j}\left(u_{V^{j 2}}, u_{V^{j 2}}\right)+3 a_{j}\left(u_{E^{j}}, u_{E^{j}}\right)
$$

For the two vertex components, we can proceed as before and find similar bounds to (31). For the edge component, we use Lemma 7.1 , the definition of $\|\cdot\|_{h, E^{j}}$,
and the fact that $\delta_{i}$ is constant at all the nodal points in $E_{h}$. We find

$$
a_{j}\left(u_{E^{j}}, u_{E^{j}}\right)=\rho_{j}\left|u_{E^{j}}\right|_{1, \Omega_{j}}^{2} \leq C \frac{\rho_{j}}{(1-\sigma)^{2}}\left\|u_{E^{j}}\right\|_{h, E^{j}}^{2} \leq C \frac{\rho_{j} \delta_{i, E}^{-2}}{(1-\sigma)^{2}}\left\|\delta_{i} u_{E^{j}}\right\|_{h, E^{j}}^{2}
$$

where $\delta_{i, E}$ is the constant value of $\delta_{i}$ on $E$. Thanks to Lemma 8.2, the norm $\|\cdot\|_{h, E^{j}}$ associated to $\Omega_{j}$ can be bounded by $\|\cdot\|_{h, E^{i}}$. In addition, we can apply Lemmas 7.4 and 7.2 and find

$$
\begin{aligned}
& \rho_{i}\left\|\delta_{i} u_{E^{j}}\right\|_{h, E^{j}}^{2} \leq C(1-\sigma)^{-2} \rho_{i}\left\|\left(\delta_{i} u_{i}\right)_{E^{i}}\right\|_{h, E^{i}}^{2} \leq C(1-\sigma)^{-2} \rho_{i}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{h, E^{i}}^{2} \\
& \quad \leq C(1-\sigma)^{-4}(1+\log k) \rho_{i}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{1, \Omega_{i}}^{2} \\
& \quad=C(1-\sigma)^{-4}(1+\log k)\left(a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u_{i}\right), \mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right)+\rho_{i} H_{i}^{-2}\left\|\mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right\|_{0, \Omega_{i}}^{2}\right)
\end{aligned}
$$

As before, the $L^{2}$ component in the last term can be bounded by the local bilinear form $a_{i}(\cdot, \cdot)$, thanks to a Poincaré or a Friedrichs inequality. Combining these two estimates and using (30), we find

$$
\begin{equation*}
a_{j}\left(u_{E^{j}}, u_{E^{j}}\right) \leq C(1-\sigma)^{-6}(1+\log k) a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u_{i}\right), \mathcal{H}_{i}\left(\delta_{i} u_{i}\right)\right) \tag{32}
\end{equation*}
$$

We next consider a substructure $\Omega_{j}$ that shares a face $F$ and thus also the edges and vertices that lie on $\partial F$. We note that on $\Omega_{j}, u$ can be decomposed as

$$
u=u_{W^{j}}+u_{F} .
$$

We have

$$
a_{j}(u, u)=\rho_{j}|u|_{1, \Omega_{j}}^{2} \leq 2 \rho_{j}\left(\left|u_{W^{j}}\right|_{1, \Omega_{j}}^{2}+\left|u_{F}\right|_{1, \Omega_{j}}^{2}\right)
$$

The wirebasket component can be bounded as before; cf. (31) and (32). For the face component we first note that the function $\delta_{i}$ is equal to a constant value $\delta_{i, F}$ at all nodal points inside $F$. Using (30), we can then write

$$
\rho_{j}\left|u_{F}\right|_{1, \Omega_{j}}^{2}=\rho_{j} \delta_{i, F}^{-2}\left|\mathcal{H}_{j}\left(\delta_{i} u_{F}\right)\right|_{1, \Omega_{j}}^{2} \leq \rho_{i}\left|\mathcal{H}_{j}\left(\delta_{i} u_{F}\right)\right|_{1, \Omega_{j}}^{2}
$$

Using Corollary 8.1 and Lemma 7.6 yields
$\left|\mathcal{H}_{j}\left(\delta_{i} u_{F}\right)\right|_{1, \Omega_{j}}^{2} \leq C(1-\sigma)^{-1}\left|\mathcal{H}_{i}\left(\delta_{i} u_{F}\right)\right|_{1, \Omega_{i}}^{2} \leq C(1-\sigma)^{-5}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}\|u\|_{1, \Omega_{i}}^{2}$.
Combining the last two estimates and using a Poincaré or a Friedrichs inequality, we find

$$
\begin{equation*}
a_{j}\left(u_{F}, u_{F}\right) \leq C(1-\sigma)^{-5}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2} a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u\right), \mathcal{H}_{i}\left(\delta_{i} u\right)\right) \tag{33}
\end{equation*}
$$

We finally need to consider the energy of $u$ in $\Omega_{i}, a_{i}(u, u)$. We note that we can decompose $u$ on $\Omega_{i}$ according to (20). The wirebasket and the face components
can be bounded as before. Summing over $i$ and the neighboring subdomains, we then find
$a(u, u) \leq \frac{C}{(1-\sigma)^{6}}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}\left(\sum_{V^{i j}} 1+\sum_{E^{i j}} 1+\sum_{F^{i j}} 1\right) a_{i}\left(\mathcal{H}_{i}\left(\delta_{i} u\right), \mathcal{H}_{i}\left(\delta_{i} u\right)\right)$.
Since the partition $\mathcal{T}_{m}$ is shape-regular, the number of subdomains to which an edge or a vertex may belong is bounded. We finally obtain

$$
\omega \leq C(1-\sigma)^{-6}\left(1+\log \left(\frac{k}{1-\sigma}\right)\right)^{2}
$$

Since in practice $\sigma$ is bounded away from one, we obtain the same bound as for Neumann-Neumann methods for $p$ finite element approximations on shaperegular meshes

$$
\kappa\left(P_{N N}\right) \leq C(1+\log k)^{2}
$$

see, e.g., [30]. We stress the fact that the constants in the last two estimates are independent of the coefficients $\rho_{i}$ and the refinement level $n$ (and thus of the aspect ratio of the mesh $\mathcal{T}_{b l}^{n, \sigma}$ ).

## 10 Numerical results

The purpose of this section is to present some numerical experiments in order to validate our analysis on some medium-size problems. An extensive numerical study is presented elsewhere; see [41].

We consider approximations on the unit cube $\Omega=(0,1)^{3}$. We choose $\rho \equiv 1$ and the right-hand side $f \equiv 1$. The macromesh $\mathcal{T}_{m}$ consists of $N \times N \times N$ cubic substructures. Geometric refinement is performed towards the three edges $x=0, y=0$, and $z=0$, with $\sigma=0.5$; see Figure 5 , left. Given a polynomial degree $k$, we choose $n=k$ as is required for robust exponential convergence; see, e.g., $[4,5]$. The conjugate gradient iteration is stopped after a reduction of the Euclidean norm of the initial residual of $10^{-14}$.

We note that even for moderate values of $k$ and $N$, extremely large linear systems are obtained; cf. Tables 1 and 2. Huge local blocks need to be inverted, both for the application of $S$ (solution of local Dirichlet problems) and the preconditioner (solution of local Neumann problems). Due to memory limitations in our Matlab implementation, we have employed approximate solvers for local Dirichlet and Neumann problems. We refer to [39, Sect. 4.4] for details on the implementation. In particular, we have used a conjugate gradient iteration with an incomplete Cholesky factorization with drop tolerance $10^{-3}$ for all local problems. The iteration is stopped after a reduction of the initial residual of a factor $10^{-3}$ or after 20 iteration steps. Our numerical results show that the theoretical bounds for the case of exact solvers in Lemma 6.2 remain valid in this case; cf. Tables 1 and 2.


Figure 5: Geometric refinement towards one corner $(N=3, \sigma=0.5$, and $n=6)$, left, and estimated condition numbers (circles) from Table 1 and least-square second order logarithmic polynomial fit (solid line) versus $k$, right.

For a fixed partition into substructures with $N=3$, Table 1 shows the size of the original problem, the iteration count, the estimated maximum and minimum eigenvalues, and the condition number for different values of $k$. We note that the minimum eigenvalue is close to one; see Lemma 6.2. In addition a moderate growth of the maximum eigenvalue is observed with $k$; such growth is consistent with the quadratic bound in Lemma 6.3; see Figure 5, right.

We next consider the same problem, and fix the polynomial degree $k=4$. Table 2 shows the results for different values of $N$. The iteration counts, and the smallest and largest eigenvalues appear to be bounded independently of the number of subdomains.

## 11 Concluding remarks

As for the analysis in [40], some important issues still need to be addressed. We refer to our previous work for a full discussion of these issues.

Our analysis is restricted to approximations that employ nodal basis functions on the Gauss-Lobatto nodes. Indeed, for three-dimensional shape-regular meshes good performance of iterative substructuring methods is in general ensured only if these basis functions are employed and for more general $p$ or $h p$ version finite element approximations many important issues remain to be solved even for shape-regular meshes; see, e.g., [38] and the references therein.

The Dirichlet and Neumann problems that we need to solve (cf. $S_{i}$ and $S_{i}^{\dagger}$ ) can be potentially very large. Approximate local solvers can be employed

| Fixed number of subdomains $(N=3)$ |  |  |  |  |  |
| :---: | :--- | :---: | :---: | :---: | :---: |
| NN (inexact) |  |  |  |  |  |
| $k$ | size | It | $\lambda_{\max }$ | $\lambda_{\min }$ | $\kappa$ |
| 2 | 1331 | 15 | 1.8379 | 1 | 1.8379 |
| 3 | 6859 | 20 | 2.8165 | 0.99997 | 2.8166 |
| 4 | 24389 | 25 | 3.9507 | 0.99947 | 3.9528 |
| 5 | 68921 | 29 | 5.1507 | 0.99799 | 5.1611 |
| 6 | 166375 | 34 | 6.3675 | 0.99801 | 6.3803 |
| 7 | 357911 | 38 | 7.5082 | 0.99395 | 7.554 |
| 8 | 704969 | 40 | 8.5298 | 0.99574 | 8.5663 |

Table 1: Conjugate Gradient method for the Schur complement system with Neumann-Neumann preconditioner with inexact solvers: iteration counts, maximum and minimum eigenvalues, and condition numbers, versus the polynomial degree, for the case of a fixed partition. The size of the original problem is also reported.
for iterative substructuring methods (see, e.g., $[39,16]$ ) and some have been proposed in [18] for $h p$ approximations. In our numerical experiments, we have employed a conjugate gradient iteration with an incomplete Choleski preconditioner, however, we believe that the tensor product structure of corner, edge, and face patches can be exploited. This is left to a future work.

We believe that the analysis and/or the development of iterative substructuring methods for general meshes with hanging nodes still need to be fully addressed. These meshes are widely used in practice. There is no straightforward way of defining Neumann-Neumann or FETI algorithms when hanging nodes lie on the interface $\Gamma$; see [40, Rem. 6.1] for more details.

Finally, our analysis has been carried out for the model problem (1), which indeed does not exhibit boundary layers. As for two-dimensional problems in [40, 42], numerical results show that our algorithms are robust when applied to certain singularly perturbed problems. Extensive numerical results will be presented in [41].

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| Fixed spectral degree $k=4$ |  |  |  |  |  |
| :---: | :--- | :---: | :---: | :---: | :---: |
| NN (inexact) |  |  |  |  |  |
| $N$ | size | It | $\lambda_{\max }$ | $\lambda_{\min }$ | $\kappa$ |
| 2 | 15625 | 18 | 2.6417 | 0.99929 | 2.6436 |
| 3 | 24389 | 25 | 3.9507 | 0.99947 | 3.9528 |
| 4 | 35937 | 28 | 4.1084 | 0.99934 | 4.1111 |
| 5 | 50653 | 29 | 4.1378 | 0.9994 | 4.1402 |
| 6 | 68921 | 30 | 4.1492 | 0.99945 | 4.1515 |
| 7 | 91125 | 30 | 4.1555 | 0.99952 | 4.1575 |
| 8 | 117649 | 30 | 4.1593 | 0.99955 | 4.1612 |
| 9 | 148877 | 30 | 4.1618 | 0.99962 | 4.1634 |
| 10 | 185193 | 30 | 4.1636 | 0.9997 | 4.1648 |

Table 2: Conjugate Gradient method for the Schur complement system with Neumann-Neumann preconditioner with inexact solvers: iteration counts, maximum and minimum eigenvalues, and condition numbers, versus the number of substructures, for the case of a fixed polynomial degree and partitions into $N \times N \times N$ substructures. The size of the original problem is also reported.

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