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Uniformly stable preconditioned mixed boundary element method for low-frequency electromagnetic scattering

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Abstract

We propose a mixed boundary element discretization of the Electric Field Integral Equation for which we have an Inf-Sup condition which is uniform in both the mesh-width h and the wave-number k, for small enough h and k. For this equation we construct a preconditioner such that the spectral condition number of the preconditioned system is also bounded independently of k and h.

The continuous problem 1

Let Ω_{-} be a bounded domain in \mathbb{R}^{3} with a smooth boundary Γ . The exterior domain $\mathbb{R}^{3} \setminus (\Omega_{-} \cup \Gamma)$ is denoted Ω_+ and the outward normal on Γ is denoted n. The tangential trace operator is denoted $\gamma_{\rm T}$ and the normal trace operator is denoted $\gamma_{\rm n}$.

Let Z be a positive constant, called impedance. For each wavenumber k > 0 the timeharmonic Maxwell equations (in any given open region of \mathbb{R}^3) are:

$$\operatorname{curl} E = +ikZ H, \quad \operatorname{curl} H = -ik/Z E.$$
(1)

Given a family (E_k^{inc}, H_k^{inc}) for small positive k of solutions of Maxwell's equations on a neighborhood of Γ we are interested, for each k, in the solution (E_k, H_k) of Maxwell's equations in Ω_{-} or Ω_{+} satisfying the perfect conductor boundary condition $\gamma_{\rm T} E_k = -\gamma_{\rm T} E_k^{inc}$, and (in the exterior domain) the Silver-Müller radiation condition.

We use potentials to represent E_k . Let G_k denote the standard Green kernel of $-\Delta - k^2$ and let Φ_k be the single layer potential defined on scalar or tangent fields u on Γ by:

$$(\Phi_k u)(y) = \int_{\Gamma} G_k(x, y) u(x) dx, \quad G_k(x, y) = \frac{e^{ik|x-y|}}{4\pi |x-y|}.$$
(2)

We represent E_k as an electric field generated by a tangent field u_k on Γ (the electric current). More precisely we put $E_k(y) = (\operatorname{grad} \operatorname{div} + k^2)(\Phi_k u_k)$. Letting $A_k = -\gamma_{\mathrm{T}}(\operatorname{grad} \operatorname{div} + k^2)\Phi_k$, the problem is to solve the Electric Field Integral Equation (EFIE) $A_k u_k = \gamma_{\rm T} E_k^{inc}$. The operator A_k is continuous from $X = {\rm H}_{\rm div}^{-1/2}(\Gamma)$ to its dual $X' = {\rm H}_{\rm rot}^{-1/2}(\Gamma)$, and the EFIE

can be put in variational form:

$$u_k \in X, \ \forall u' \in X \quad \langle A_k u_k, u' \rangle = \langle E_k^{inc}, u' \rangle. \tag{3}$$

We denote by a_k the associated bilinear form ; its expression on smooth fields is:

$$a_k(u,v) = \iint_{\Gamma \times \Gamma} G_k(x,y) (\operatorname{div} u(x) \operatorname{div} v(y) - k^2 u(x) \cdot v(y)) \mathrm{d}x \mathrm{d}y \tag{4}$$

Following Bendali [1] this variational problem is solved with the Galerkin method on divconforming Finite Element spaces on the boundary. At low frequencies one sees that the problem is that the limit of the operator A_k as $k \to 0$, is degenerated; in fact the limit is not even Fredholm since its kernel contains the infinite dimensional space rot $\mathrm{H}^{1/2}(\Gamma)$.

The object of this paper is to compute approximations of u_k in a stable way for small k.

2 The continuous remedy

For simplicity we suppose that Γ is connected and simply connected. As remarked by De-LaBourdonnaye [2], if we put $V = \operatorname{grad} H^{3/2}(\Gamma)$ and $W = \operatorname{rot} H^{1/2}(\Gamma)$, then V and W are closed in X and we have the decomposition:

$$X = V \oplus W. \tag{5}$$

We put $S = H^{1/2}(\Gamma)$, and for any space Y of scalar fields on Γ we put $Y^{\bullet} = \{u \in Y : \langle u, 1 \rangle = 0\}$.

Let $\Xi_k : V \times S^{\bullet} \to X$ denote the isomorphism defined by $\Xi_k(v, p) = v + k^{-1}$ rot p. The four blocks of the bilinear form \tilde{a}_k on $V \times S^{\bullet}$ defined by $\tilde{a}_k((v, p), (v', p')) = a_k(\Xi_k(v, p), \Xi_k(v', p'))$, have the expression:

$$\begin{pmatrix}
\iint G_k(x,y)(\operatorname{div} v(x) \operatorname{div} v'(y) - k^2 v(x) \cdot v'(y)) \mathrm{d}x \mathrm{d}y & -k \iint G_k(x,y) v(x) \cdot \operatorname{rot} p'(y) \mathrm{d}x \mathrm{d}y \\
-k \iint G_k(x,y) \operatorname{rot} p(x) \cdot v'(y) \mathrm{d}x \mathrm{d}y & -\iint G_k(x,y) \operatorname{rot} p(x) \cdot \operatorname{rot} p'(y) \mathrm{d}x \mathrm{d}y \\
\end{cases}$$
(6)

Since there is C > 0 such that:

$$\forall v \in V \quad \|v\|_X \le C \|\operatorname{div} v\|_{\mathrm{H}^{-1/2}(\Gamma)}, \quad \forall p \in S^{\bullet} \quad \|p\|_S \le C \|\operatorname{rot} p\|_{\mathrm{H}^{-1/2}_{\mathrm{T}}(\Gamma)}, \tag{7}$$

the two diagonal blocks are coercive hence invertible for k = 0. We remark also that the coupling blocks vanish for k = 0.

Concerning the right-hand sides we remark that:

$$k^{-1} \langle \gamma_{\rm T} E_k^{inc}, \operatorname{rot} p' \rangle = i Z \langle \gamma_{\rm n} H_k^{inc}, p' \rangle.$$
 (8)

If for instance the family of incident waves consists of plane waves:

$$E_k^{inc}(x) = E_0 e^{ik\sigma \cdot x}, \quad H_k^{inc}(x) = 1/ZE_0 \times \sigma e^{ik\sigma \cdot x}, \tag{9}$$

then the limit of $\gamma_{\rm T} E_k^{inc}$ is a surface gradient and $\gamma_{\rm n} H_k^{inc}$ has a non-zero limit. Thus both righthand sides in $V^* \times S^{\bullet \star}$ have non-zero limits as $k \to 0$. It follows that with the decomposition $u_k = \Xi_k(v_k, p_k)$ both v_k and p_k have a non-zero limit as $k \to 0$.

We now turn to the preconditioning of the variational problem associated with (6) and we recall the remark made in [3] that a preconditioner is obtained by an invertible bilinear form on a dual space. Since the off-diagonal terms are small in norm and compact it is enough to precondition the two diagonal blocks.

For the first block, we proceed as follows. Put $V' = \operatorname{grad} \operatorname{H}^{1/2}(\Gamma)$. Then we remark that the $\operatorname{L}^2_{\operatorname{T}}(\Gamma)$ -bilinear form extends continuously to an invertible bilinear form on $V' \times V$. Let $\Theta_1 : V^* \to V'$ be the corresponding isomorphism. We remark furthermore that V' is a closed subspace of $\operatorname{H}^{-1/2}_{\operatorname{T}}(\Gamma)$, hence we can use the bilinear form associated with the single layer operator on tangent fields as a preconditioner.

For the second block, the induced operator on $S^{\bullet} \to S^{\bullet \star}$ is the main part of the hypersingular operator appearing in acoustics. It can be efficiently preconditioned by the single layer operator [6] [3]. As a matter of notations we put $S' = \mathrm{H}^{-1/2}(\Gamma)$ so that the $\mathrm{L}^2(\Gamma)$ -bilinear form extends continuously to an invertible bilinear form on $S'^{\bullet} \times S^{\bullet}$, and let $\Theta_2 : S^{\bullet \star} \to S'^{\bullet}$ be the corresponding isomorphism.

Thus, on $V' \times S'^{\bullet}$ we use the bilinear form b whose block expression is:

$$\begin{pmatrix}
\iint G_0(x,y)v(x) \cdot v'(y) \mathrm{d}x \mathrm{d}y & 0 \\
0 & -\iint G_0(x,y)q(x)q'(y) \mathrm{d}x \mathrm{d}y
\end{pmatrix}$$
(10)

Letting $\Theta : V^* \times S^{\bullet *} \to V' \times S'^{\bullet}$ be the map componentwise induced by Θ_1 and Θ_2 , and associating an operator $\tilde{\mathcal{A}}_k : V \times S^{\bullet} \to V^* \times S^{\bullet *}$ with \tilde{a}_k and $\mathcal{B} : V' \times S'^{\bullet} \to V'^* \times S'^{\bullet *}$ with b we have:

Proposition 2.1 There is $\epsilon > 0$ such that for all $k \in [0, \epsilon]$ the operator $\Theta^* \mathcal{B} \Theta \mathcal{A}_k$ is an automorphism of $V \times S^{\bullet}$ and all terms of the composition are isomorphisms whose norm and norm of the inverse are bounded independently of k in $[0, \epsilon]$.

3 Discretization

Since $V = \text{grad } \mathrm{H}^{3/2}(\Gamma)$ it would be cumbersome to implement a conforming Finite Element discretization of the variational problem on $V \times S^{\bullet}$. Instead we propose the following non-conforming method.

Suppose we have (finite dimensional) subspaces X_h of $X \cap H^0_{div}(\Gamma)$ and S_h of $S \cap H^1(\Gamma)$, which are stable under complex conjugation, which are such that S_h contains the constant fields and we have an exact sequence:

$$S_h \xrightarrow{\text{rot}} X_h \xrightarrow{\text{div}} L^2(\Gamma).$$
 (11)

We define V_h by:

$$V_h = \{ u \in X_h : \forall p \in S_h \quad \langle u, \operatorname{rot} p \rangle = 0 \}.$$
(12)

We keep the notation \tilde{a}_k to denote the extension of \tilde{a}_k to $X \times S$ whose block-wise expression is given by (6). We solve the system: Find $(v_{kh}, p_{kh}) \in V_h \times S_h^{\bullet}$, such that for all $(v', p') \in V_h \times S_h^{\bullet}$ we have:

$$\tilde{a}_k((v_{kh}, p_{kh}), (v', p')) = \langle E_k^{inc}, v' \rangle + iZ \langle H_k^{inc} \cdot n, p' \rangle.$$
(13)

Recall the definition of the gap : $\delta_X(V_h, V) = \sup_{v_h \in V_h} \inf_{v \in V} ||v_h - v||_X / ||v_h||_X$. Our first proposition concerns the well posedness of the discrete system.

Proposition 3.1 If $\delta_X(V_h, V) \to 0$ as $h \to 0$ then there is $\epsilon > 0$, $h_0 > 0$ and C > 0 such that for all $k \in [0, \epsilon]$, all $h < h_0$ we have :

$$\inf_{(v,p)\in V_h\times S_h^{\bullet}} \sup_{(v',p')\in V_h\times S_h^{\bullet}} \frac{|\tilde{a}_k((v,p),(v',p'))|}{\|(v,p)\|_{X\times S}\|(v',p')\|_{X\times S}} \ge 1/C.$$
(14)

-*Proof:* Actually we prove uniform coercivity. By the continuity of the operators with respect to k it suffices to prove coercivity for k = 0. Let P be the projector with range V and kernel W. We have:

$$\forall v \in V_h \ \|v\|_X \leq \|v - Pv\|_X + \|Pv\|_X \tag{15}$$

$$\leq \|I - P\|\delta(V_h, V)\|v\| + \|\operatorname{div} v\|_{\mathrm{H}^{-1/2}(\Gamma)}.$$
 (16)

Therefore we have an estimate of the form : There is $h_0 > 0$ and C > 0 such that for all $h < h_0$:

$$\forall v \in V_h \quad ||v||_X \le C ||\operatorname{div} v||_{\mathrm{H}^{-1/2}(\Gamma)}.$$
 (17)

The result entails.

In general we do not have a basis of V_h , hence solving this system requires some extra work. In our case this will be carried out by the preconditioner which we define now. It should be checked that in what follows only bases of X_h and S_h are needed.

Let $\Theta_{1h}: X_h^* \to X_h$ denote the map which to any $\ell \in X_h^*$ associates the solution of:

$$v \in V_h, \ \forall v' \in V_h \quad \langle v, v' \rangle = \ell(v').$$
 (18)

For $\ell \in X_h^{\star}$, $\Theta_{1h}\ell$ can be computed simply as the solution u of:

$$p \in S_h^{\bullet}, \ \forall p' \in S_h^{\bullet} \quad \langle \operatorname{rot} p, \operatorname{rot} p' \rangle = \ell(\operatorname{rot} p'), \\ u \in X_h, \ \forall u' \in X_h \quad \langle u, u' \rangle = \ell(u') - \langle \operatorname{rot} p, u' \rangle.$$

$$(19)$$

We define the discretization Θ_{2h} of Θ_2 to be the map which to $\ell \in S_h^*$ associates the solution p of:

$$p \in S_h^{\bullet}, \ \forall p' \in S_h^{\bullet} \quad \langle p, p' \rangle = \ell(p').$$
 (20)

Let $\Theta_h : X_h^{\star} \times S_h^{\star} \to X_h \times S_h$ be the association of Θ_{1h} and Θ_{2h} . We keep the notation bfor the extension of b from $V' \times S'^{\bullet}$ to $\mathrm{H}_{\mathrm{T}}^{-1/2}(\Gamma) \times S'$ keeping the block expression (10). Let $\mathcal{B}_h : X_h \times S_h \to X_h^{\star} \times S_h^{\star}$ be the map induced by b. We also denote by $\tilde{\mathcal{A}}_{kh}$ the map induced by \tilde{a}_k on $X_h \times S_h \to X_h^{\star} \times S_h^{\star}$.

One sees that the operator $\Theta_h^* \mathcal{B}_h \Theta_h$ is a surjection onto $V_h \times S_h^{\bullet}$. For $\ell \in (X_h \times S_h^{\bullet})^*$, $\Theta_h^* \mathcal{B}_h \Theta_h \ell$ depends only on $\ell|_{V_h \times S_h^{\bullet}}$. It follows that the conjugate gradient algorithm for $\tilde{\mathcal{A}}_{kh}$ on $X_h \times S_h^{\bullet}$, preconditioned by $\Theta_h^* \mathcal{B}_h \Theta_h$ yields iterates in $V_h \times S_h^{\bullet}$ converging to the solution of (13). Moreover $\Theta_h^* \mathcal{B}_h \Theta_h \tilde{\mathcal{A}}_{kh}$ determines a bijection $V_h \times S_h^{\bullet} \to V_h \times S_h^{\bullet}$ whose spectral condition number κ_{kh} is bounded independently of k in an interval $[0, \epsilon]$.

More precise estimates on κ_{kh} and the convergence of Krylov subspace methods, depend on the actual Galerkin spaces. Examples of Finite Element spaces which satisfy the above conditions include the case where we have quasi-uniform triangulations of Γ and take for X_h Raviart-Thomas vector FE of degree n and for S_h the scalar continuous piecewise P^{n+1} FE. Then we also have the following stability property:

Proposition 3.2 There is $\epsilon > 0$, $h_0 > 0$ and $C_1, C_2 > 0$ such that for all $k \in [0, \epsilon]$, all $h < h_0$ we have:

$$\begin{aligned} \|\Theta_{h}^{\star} \mathcal{B}_{h} \Theta_{h} \tilde{\mathcal{A}}_{kh}(u, p)\|_{0} &\leq C_{1} \|(u, p)\|_{0} \\ & \text{with } \|(u, p)\|_{0}^{2} = \|u\|_{\mathrm{H}^{0}_{\mathrm{div}}(\Gamma)}^{2} + \|p\|_{\mathrm{H}^{1}(\Gamma)}^{2}, \end{aligned}$$
(21)

$$\begin{aligned} \|\Theta_{h}^{\star}\mathcal{B}_{h}\Theta_{h}\tilde{\mathcal{A}}_{kh}(u,p)\|_{-1/2} &\geq C_{2}^{-1}\|(u,p)\|_{-1/2} \\ \text{with } \|(u,p)\|_{-1/2}^{2} = \|u\|_{\mathrm{H}^{-1/2}_{\mathrm{div}}(\Gamma)}^{2} + \|p\|_{\mathrm{H}^{1/2}(\Gamma)}^{2}. \end{aligned}$$
(22)

-Proof: We show the proof for the case of fixed k = 0 and then indicate how the proof extends to the the case of k in an interval $[0, \epsilon]$.

At k = 0 the operator \mathcal{A}_{kh} decouples and we can study the action on V_h and S_h^{\bullet} separately.

(i) Proof of estimate (21). Let P_{V_h} denote the $L^2_T(\Gamma)$ -orthogonal projection onto V_h , and $P_{S_h^{\bullet}}$ denote the $L^2(\Gamma)$ -orthogonal projection onto S_h^{\bullet} . Let \mathfrak{S} and \mathfrak{S}_T denote the single layer operator acting on scalar and tangent fields respectively.

The discrete operator $\Theta_h^* \mathcal{B}_h \Theta_h \tilde{\mathcal{A}}_{kh}$ decouples into the two operators:

$$V_h \subset \mathrm{H}^0_{\mathrm{div}}(\Gamma) \xrightarrow{\mathrm{div}} \mathrm{L}^2(\Gamma) \xrightarrow{\mathfrak{S}} \mathrm{H}^1(\Gamma) \xrightarrow{\mathrm{grad}} \mathrm{L}^2_{\mathrm{T}}(\Gamma) \xrightarrow{P_{V_h}} V_h \xrightarrow{\mathfrak{S}_{\mathrm{T}}} \mathrm{H}^1_{\mathrm{T}}(\Gamma) \xrightarrow{P_{V_h}} V_h, \qquad (23)$$

and:

$$S_{h}^{\bullet} \subset \mathrm{H}^{1}(\Gamma) \xrightarrow{\mathrm{rot}} \mathrm{L}_{\mathrm{T}}^{2}(\Gamma) \xrightarrow{\mathfrak{S}_{\mathrm{T}}} \mathrm{H}_{\mathrm{T}}^{1}(\Gamma) \xrightarrow{\mathrm{rot}} \mathrm{L}^{2}(\Gamma) \xrightarrow{P_{S_{h}^{\bullet}}} S_{h}^{\bullet} \xrightarrow{\mathfrak{S}} \mathrm{H}^{1}(\Gamma) \xrightarrow{P_{S_{h}^{\bullet}}} S_{h}^{\bullet}.$$
(24)

For the first operator it thus suffices to show that the last occurrence of P_{V_h} satisfies an estimate of the form:

$$\forall u \in \mathrm{H}^{1}_{\mathrm{T}}(\Gamma) \quad \|P_{V_{h}}u\|_{\mathrm{H}^{0}_{\mathrm{div}}(\Gamma)} \leq C\|u\|_{\mathrm{H}^{1}_{\mathrm{T}}(\Gamma)}.$$

$$(25)$$

Let P_{X_h} denote the L^2_T -orthogonal projection onto X_h and let Π_h denote the standard interpolator onto X_h (interpolating the fluxes through the edges of the curved triangles). Using the fact that V_h is L^2_T -orthogonal to the kernel of div on X_h , and standard inverse inequalities we have for any $u \in H^1_T(\Gamma)$:

$$\|\operatorname{div} P_{V_h} u\|_{L^2(\Gamma)} = \|\operatorname{div} P_{X_h} u\|_{L^2(\Gamma)}$$
(26)

$$\leq \| \operatorname{div}(P_{X_{h}}u - \Pi_{h}u) \|_{\mathrm{L}^{2}(\Gamma)} + \| \operatorname{div}\Pi_{h}u \|_{\mathrm{L}^{2}(\Gamma)}$$
(27)

$$\leq Ch^{-1} \|\operatorname{div}(P_{X_h}u - \Pi_h u)\|_{\mathrm{H}^{-1}(\Gamma)} + \|\operatorname{div}\Pi_h u\|_{\mathrm{L}^2(\Gamma)}$$
(28)

$$\leq Ch^{-1} \| P_{X_h} u - \Pi_h u \|_{\mathrm{L}^2_{\mathrm{T}}(\Gamma)} + \| \operatorname{div} \Pi_h u \|_{\mathrm{L}^2(\Gamma)}$$
(29)

$$\leq Ch^{-1} \| u - \Pi_h u \|_{L^2_{T}(\Gamma)} + \| \operatorname{div} \Pi_h u \|_{L^2(\Gamma)}$$
(30)

The estimate (25) now follows from the well-known properties of Π_h .

For the second operator the $\mathrm{H}^1(\Gamma)$ -stability follows from the well-known $\mathrm{H}^1(\Gamma)$ -stability of the $\mathrm{L}^2(\Gamma)$ -projector onto S_h^{\bullet} (in the last occurrence of $P_{S_h^{\bullet}}$).

Thus estimate (21) is proved.

(ii) Proof of estimate (22). Since \mathcal{B}_h and $\tilde{\mathcal{A}}_{0h}$ are coercive on $V_h \times S_h^{\bullet}$ in the $\mathrm{H}_{\mathrm{T}}^{-1/2}(\Gamma) \times \mathrm{H}^{-1/2}(\Gamma)$ and $\mathrm{H}_{\mathrm{div}}^{-1/2}(\Gamma) \times \mathrm{H}^{1/2}(\Gamma) \times \mathrm{H}^{1/2}(\Gamma)$ norms respectively it suffices to show that the L²-dualities used for the projections are uniformly (with respect to h) continuous on V_h and S_h^{\bullet} in these norms. For the L²-duality on S_h in the $\mathrm{H}^{-1/2}(\Gamma) \times \mathrm{H}^{1/2}(\Gamma)$ norm this is trivial. For the case of V_h we proceed as follows:

Let P denote the projector with range V and kernel W. It preserves the divergence. As we remarked in [4] we have an estimate of the form:

$$\forall v \in V_h \quad \|v - Pv\|_{\mathcal{L}^2_{\mathcal{T}}(\Gamma)} \le Ch\| \operatorname{div} v\|_{\mathcal{L}^2(\Gamma)}.$$
(31)

For any $v, v' \in V_h$ we have:

$$|\langle v, v' \rangle| \leq |\langle v - Pv, v' \rangle| + |\langle Pv, v' \rangle|$$
(32)

$$\leq C \|v - Pv\|_{\mathrm{H}^{0}_{\mathrm{div}}(\Gamma)} \|v'\|_{\mathrm{H}^{-1}_{\mathrm{rot}}(\Gamma)} + \|Pv\|_{\mathrm{H}^{1/2}_{\mathrm{T}}} \|v'\|_{\mathrm{H}^{-1/2}_{\mathrm{T}}(\Gamma)}$$
(33)

$$\leq Ch \|\operatorname{div} v\|_{\mathrm{L}^{2}(\Gamma)} \|v'\|_{\mathrm{L}^{2}_{\mathrm{T}}(\Gamma)} + C\|\operatorname{div} v\|_{\mathrm{H}^{-1/2}(\Gamma)} \|v'\|_{\mathrm{H}^{-1/2}_{\mathrm{T}}(\Gamma)}$$
(34)

$$\leq C \|\operatorname{div} v\|_{\mathrm{H}^{-1/2}(\Gamma)} \| \|v'\|_{\mathrm{H}^{-1/2}_{\pi}(\Gamma)}.$$
(35)

This completes the proof of estimate (22).

To extend the results to to $k \in [0, \epsilon]$ we simply remark that \mathcal{B}_h and \mathcal{A}_{kh} are uniformly continuous for the norms used in the proof of (21) and uniformly coercive for the norms used in the proof of (22), with respect to k in an interval $[0, \epsilon]$.

It follows that κ_{kh} is bounded by C_1C_2 for (k, h) in the range $[0, \epsilon] \times]0, h_0[$.

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